

Solid-state formation of Titanium Carbide and Molybdenum Carbide as contacts for carbon-containing semiconductors

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Abstract

Metal carbides are good candidates to contact carbon-based semiconductors (SiC, diamond, carbon nanotubes). Here, we report on an *in-situ* study of carbide formation during the solid-state reaction between thin Ti or Mo films and C substrates. Titanium Carbide, TiC , was previously reported as a contact material to diamond and carbon nanotubes. However, the present study shows two disadvantages for the solid-state reaction of Ti and C. First, because Ti reacts readily with oxygen, a capping layer should be included to enable carbide formation. Second, the TiC phase can exist over a wide range of composition (about 10%, i.e. from $Ti_{0.5}C_{0.5}$ to $Ti_{0.6}C_{0.4}$), leading to significant variations in the properties of the material formed. The study of the Mo-C system suggests that Molybdenum Carbide, Mo_2C , is a promising alternative, since the phase shows a lower resistivity (about 45% lower than for TiC), the carbide forms below $900^\circ C$, and its formation is less sensitive to oxidation as compared to the Ti-C system. The measured resistivity for Mo_2C is $\rho = 59\mu\Omega cm$, and from kinetics studies an activation energy for Mo_2C formation of $E_a = 3.15 \pm 0.15 eV$ was obtained.

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I. INTRODUCTION

Several carbon-based semiconductors are promising materials for future electronic applications. SiC offers significant advantages for high power-switching devices¹. Doped diamond is a superior wide-band-gap semiconductor, offering a high breakdown voltage, high thermal conductivity, small dielectric constant and excellent radiation hardness². Carbon nanotubes (CNTs) are quasi-one-dimensional molecular structures with semiconducting or metallic properties, which are widely investigated for application in the field of nano-electronics³. While these advantages are being widely investigated, there is also a need for reliable contacts to these carbon-containing semiconductors.

In silicon-based technology, metal-silicon compounds (silicides) are now widely used as contact materials, since they offer several advantages as compared to metal contacts. Unlike pure metals, several silicide phases are thermodynamically stable in contact with silicon and, when silicides are formed by a solid-state reaction between metal and Si, contact formation can be achieved by means of a self-aligned process (SALICIDE)⁴. The solid-state reaction between metal and Si only occurs in those regions of a patterned substrate, where the metal is in direct contact with the Si (or C for the present case of carbides). Therefore, there is no need for alignment and lithographic patterning of the metal film, since the formation of the silicide will be automatically restricted to the contact region. One can expect the same advantages when using carbides to contact carbon-based semiconductors. While titanium has already been used to contact diamond^{5,6} and CNT's⁷, relatively little work has been reported concerning carbide formation during thin film reactions.

In this work, we studied the formation of titanium carbide and molybdenum carbide through a solid-state reaction. Ti was chosen because of the existence of prior work⁵. While the Mo is shown to react with C at higher temperature, the inactivity of Mo with respect to oxidation presents clear process advantages.

II. EXPERIMENTAL

The samples used in this work consisted of 30nm Ti or Mo on different substrates: SiO_2 , CNT's deposited on SiO_2 , or 200nm amorphous C on SiO_2 . Sometimes capping layers were included (5 or 30nm, amorphous C or TiN), and they will be mentioned where used. All

layers were sputter deposited at a base pressure of 10^{-6} mbar . The sputtering equipment was calibrated prior to sample preparation, using several depositions of each metal and a Talystep Profilometer.

Several techniques were used for this study: four point probe for resistance measurements, XPS for phase composition information, RBS for contamination and thickness measurements, and *in-situ* XRD for phase identification and for determining the activation energy E_a . The *in-situ* x-ray diffraction (XRD) experiments were performed at the X20C-beamline of the National Synchrotron Light Source (NSLS) at Brookhaven National Laboratory. During annealing, the sample is continuously illuminated by an intense beam of monochromatic X-rays ($\lambda = 0.177$ to 0.180 nm). During the measurement, the samples were heated from 100°C to 1000°C in a purified *He*-atmosphere, at different ramp rates (from 0.3°C/s to 35°C/s). Diffracted photons from the sample are continuously detected by an array of photodiodes covering 14° in 2θ . Acquisition times can vary from 100 ms to several minutes depending on the annealing conditions. For the standard heating rate of 3°C/s presented here, a spectrum is taken every 0.5 s.

III. RESULTS AND DISCUSSION

A. The Ti-C system

Figure 1 shows *in-situ* XRD patterns of 30nm Ti deposited on SiO_2 , and on 200nm amorphous carbon (with and without a 30nm carbon capping layer). The XRD data is presented both as contours and as a gray scale (black is the highest intensity). For the as-deposited sample, one can easily identify the $\text{Ti}(002)$ peak near $2\theta \approx 45^\circ$. Upon annealing to 1000°C , this peak shifts towards lower 2θ -values for samples without a C capping layer, indicating an expansion of the Ti lattice. This is caused by oxidation of the Ti-layer, as confirmed by ex-situ RBS analysis of quenched samples. This oxidation can either be a reaction with the oxygen impurities left in the annealing chamber, or with oxygen from the substrate. An XPS measurement shown in Figure 2(a), confirms that the uncapped Ti mainly oxidizes. However, analysis of the single C-peak spectra showed that there is some TiC -formation (insets in Figure 2). The C spectrum at depth (1) shows a peak at lower binding energy (marked with ' $\text{C}(\text{TiC})\text{-peak}$ '), that indicates that there is TiC -formation.

The second spectrum in the insets show the standard C1s -peak, and was obtained at depth (2), where no metal and thus no carbide formation is present.

When an extra 30nm C-layer is included as a capping layer, oxidation is prevented, and one observes the formation of a titanium carbide between $T_f^{TiC} = 550^\circ C$ and $650^\circ C$ (Figure 1(c)). There are diffraction peaks around $2\theta \approx 42^\circ$ and $\approx 49^\circ$, corresponding to the (111)-peak and the (200)-peak of TiC . No other diffraction peaks could be observed from $2\theta \approx 35^\circ$ to $\approx 70^\circ$. Because there are also other possible phases of titanium carbide with diffraction peaks at these angles of 2θ (e.g. Ti_8C_5), XPS was performed on this sample. The formation of the TiC phase could be confirmed on the basis of the Ti/C ratio in the layer (Figure 2(b)). The difference in oxygen concentration between the samples with and without a capping layer is very clear.

From the above, it is clear that the sensitivity of Ti towards oxidation is a major problem when attempting to form TiC . A second problem of TiC is related to the solubility of C. TiC can be stable in a range of about 10% of C solubility⁸. With a lower C solubility, there will be C vacancies, which will act as powerful scattering centers for electrons and phonons. This will reflect itself in certain physical properties of the material, e.g. diffusion, thermal conductivity and (most important) the resistivity^{9,10}. For instance, for the TiC layer formed in Figure 1(c), a resistivity of $\rho = 136\mu\Omega cm$ was measured using a standard four-point probe. Earlier work reports on values ranging from 35 to $250\mu\Omega cm$. Williams¹¹ measured room-temperature resistivities for different C-vacancy concentrations (from $TiC_{0.93}$ to $TiC_{0.83}$). For stoichiometric TiC , Williams obtained a room-temperature resistivity of $70 \pm 10\mu\Omega cm$. Up to 7% C-vacancies, the resistivity shows a monotonic increase with increasing vacancy concentration, with an added resistivity of $16\mu\Omega cm$ per atomic percent of carbon vacancies in TiC . At higher vacancy concentrations (which means a lower C content, i.e. TiC_x with $x \leq 0.92$), the resistivity saturates around $200\mu\Omega cm$. According to this model, our TiC films should have a C vacancy concentration of about 4%. This is probably an upper limit, because impurities (like oxygen) will also have an influence on the resistivity of the film.

A thickness for the TiC film was calculated from crystallographic data¹² $t = 34.3nm$, which is in agreement with the RBS measurements ($t = 36.7 \pm 2nm$).

B. The Mo-C system

According to metal-C binary phase diagrams⁸, there are 14 metals other than Ti that can form stable carbide-phases: Al, Sc, V, Cr, Mn, Fe, Y, Zr, Nb, Mo, Hf, Ta, W and U. When we take a closer look at the binary phase diagrams, we see that only 7 of these metal-C systems have a narrow range of composition (observed as a single vertical line on the binary phase diagrams): Al, Cr, Mn, Fe, Mo, W and U, while the others exhibit carbide phases with a wide range of C solubility (like Ti). From the compounds with a narrow range of composition, we have selected molybdenum as a good candidate, because of its relative insensitivity to oxygen. This is demonstrated in Figure 3(a), where a 30nm Mo layer on SiO_2 is annealed at $3^\circ C/s$ to $1000^\circ C$, and no reaction is observed between the metal layer and the substrate, nor between the metal layer and the annealing atmosphere. In Figure 3(b) the formation of the Mo_2C -carbide can be seen between $T_f = 830^\circ C$ and $900^\circ C$. As evidenced by the disappearance of the Mo(110)- peak around $2\theta \approx 47^\circ$, there is a clear phase transition from the metal phase to the Mo_2C phase. We can see carbide peaks around $2\theta \approx 40^\circ$, $\approx 44^\circ$ and $\approx 46^\circ$, corresponding to the $Mo_2C(100)$, $Mo_2C(002)$ and $Mo_2C(101)$ planes, respectively. Other Mo_2C diffraction peaks could be observed when positioning the detector in other 2θ -windows (not shown here). Figure 4 shows the solid-state reaction for a Mo sample with different capping layers: 5nm and 30nm C, and a 30nm TiN capping layer. The transition to Mo_2C was observed for all samples, although the formation temperature varied significantly: transition temperatures between $T_f^{5nmC} = 730^\circ C$ and $800^\circ C$, $T_f^{TiN} = 750^\circ C$ and $830^\circ C$ and $T_f^{30nmC} = 700^\circ C$ and $770^\circ C$ were observed, for the 5nm C, 30nm TiN and 30nm C capping layers respectively. The difference in formation temperature (which can be observed in Figure 4) can be attributed to a competition between C, O and N, for the interstitial sites⁹. Without a capping layer, some impurities from the annealing atmosphere can enter the Mo lattice and occupy interstitial sites. These interstitial locations facilitate C diffusion and are therefore important during carbide formation. As a result, the more interstitial sites that are occupied by other elements than C, the harder it is to form the carbide phase, and the higher the formation temperature. We observe that a 5nm C capping layer already reduces the O diffusion within the Mo film, which makes it easier for the C to occupy the interstitial sites, and we get a lower formation temperature. Increasing the thickness of the C capping layer to 30nm, blocks the O diffusion even more and thus the formation temperature for

the carbide is lowered again. When a 30nm TiN capping layer is used, there may be some N moving into the Mo lattice, which would explain the increase in formation temperature, compared to the C capping layers.

We measured the resistivity for the Mo_2C phase $\rho = 59\mu\Omega cm$, and calculated the thickness from crystallographic data¹², $t = 35.8nm$, which agrees with the thickness determined by RBS ($t = 38.4 \pm 1nm$). We also performed *in-situ* XRD measurements at heating rates of 0.3, 1, 3, 5, 9 and $35^\circ C/s$. Using the Kissinger formula¹³:

$$\ln \left(\frac{dT/dt}{T_f^2} \right) = -\frac{E_a}{k_b T_f} + C, \quad (1)$$

where dT/dt is the heating rate, T_f is the formation temperature and k_b the boltzmann constant, a value of $E_a = 3.15 \pm 0.15 eV$ was found for the activation energy for Mo_2C formation. Activation energies for C diffusion in carbides are relatively high compared to other systems (e.g. $4.5eV$ for C in TiC ¹⁴ versus $2.5eV$ for Si in $TiSi_2$, which is the largest value reported in¹⁵). Also, from literature^{16,17} one knows that the activation energy is usually much higher for C diffusion in a carbide layer, than in the metal matrix (e.g. $1.89eV$ for C diffusion in Mo). Values reported for C diffusion in Mo_2C range from $3.06eV$ ¹⁸ to $3.60eV$ ¹⁷ in different temperature ranges. Warnes and Simkovich¹⁹ reported a much higher activation energy of $4.79eV$ in a temperature range between $800^\circ C$ and $1000^\circ C$. The value for the activation energy, obtained in this work, is in the same range as reported in¹⁸ and¹⁷.

C. Reaction with CNT's and Diamond

In this work, amorphous C was used as a model system to investigate the carbide formation. To verify if the results obtained for reaction of a metal with a-C directly relate to the reaction with diamond or CNT's, the reaction of Mo was also tested on these C-substrates. Therefore, a layer of carbon nanotubes (CNT's) was dispersed on the SiO_2 -substrate and covered with 30nm of Mo. In Figure 5 the *in-situ* XRD pattern of the reaction of Mo and CNT's can be compared to the reaction of Mo and a-C, which was discussed above. With the Mo-CNT system, a part of the Mo reacts with these CNT's to form the stable Mo_2C , while the remaining Mo stays unreacted on the substrate. Notice that the Mo_2C phase is formed at about the same temperature for the a-C sample and the CNT-sample, showing that $T_f^{Mo_2C}$ is independent of the chemical nature of the substrate. The same experiment

was done with Ti and CNT's, but the Ti tends to react with the SiO_2 -substrate and/or with the oxygen impurities left in the annealing chamber, instead of reacting solely with the CNT's. This may cause problems when using Ti in advanced devices with CNT's. As further evidence, 30nm Mo was deposited on CVD diamond and annealed to $1000^\circ C$ at the standard heating rate of $3^\circ C/s$. Figure 6 shows standard $\theta/2\theta$ XRD patterns measured *ex-situ* (using $CuK\alpha$ radiation). There is a clear difference between the as-deposited sample (a) and the annealed sample (b). This evidences that the carbide formation for the Mo-C system is independent of the chemical nature of the carbon-to-carbon bonding in the substrate.

In this work, the electric nature of the carbide-carbon interface hasn't been investigated, but the mere formation of the carbide was focussed on. The use of amorphous carbon for this, didn't allow investigating the contact resistance. However, Nakanishi *et al.*²⁰ annealed Mo contacts to boron-doped polycrystalline diamond films for 60 min at $600^\circ C$, to find a low contact resistance of $10^{-6}\Omega cm^2$. They found the phase to be extremely thermally stable, and found that the Mo_2C carbide was formed. For Ti contacts, they found the optimized annealing condition to be $400^\circ C$ for 4 min. They applied an isothermal annealing (in our work, it was a ramp anneal), but the same trend in formation temperature is observed $T_f^{TiC} < T_f^{Mo_2C}$.

IV. CONCLUSION:

In conclusion, we report the formation of molybdenum carbide (Mo_2C) and titanium carbide (TiC) through a solid-state reaction. Mo_2C is a promising alternative to TiC , because its formation is less sensitive to the presence of oxygen. Secondly, the compound Mo_2C only exists over a narrow range of composition (unlike TiC), which will reflect itself in more well-defined physical properties. Also, we have shown for the Mo-C system that the formation temperature is independent of the chemical nature of the carbon-to-carbon bonding in the substrate, and that Mo can be used to form a stable carbide phase with carbon nanotubes and diamond.

Acknowledgments

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- ¹ J. A. Cooper and A. Agarwal, Proc. IEEE **90**, 956 (2002).
- ² V. K. Bazhenov, I. M. Vikulin, and A. G. Gontar, Sov. Phys. Semicond. **19**, 829 (1985).
- ³ P. Avouris, Chem. Phys. **281**, 429 (2002).
- ⁴ S. L. Zhang and M. Ostling, Crit. Rev. Solid State Mat. Sci. **28**, 1 (2003).
- ⁵ T. Tachibana, B. E. Williams, and J. T. Glass, Phys. Rev. B **45**, 11975 (1992).
- ⁶ K. L. Moazed, J. R. Zeidler, and M. J. Taylor, J. Appl. Phys. **68**, 2246 (1990).
- ⁷ R. Martel, V. Derycke, C. Lavoie, J. Appenzeller, K. K. Chan, J. Tersoff, and P. Avouris, Phys. Rev. Lett. **87** (2001).
- ⁸ T. B. Massalski, Metall. Trans. B **20**, 445 (1989).
- ⁹ H. J. Goldschmidt, *Interstitial Alloys* (Butterworths, London, 1967).
- ¹⁰ W. S. Williams, JOM-J. Miner. Met. Mater. Soc. **49**, 38 (1997).
- ¹¹ W. S. Williams, Phys. Rev. **135**, A505 (1964).
- ¹² J. R. Chen, J. Vac. Sci. Technol. A-Vac. Surf. Films **5**, 1802 (1987).
- ¹³ E. G. Colgan and F. M. dHeurle, J. Appl. Phys. **79**, 4087 (1996).
- ¹⁴ S. Sarian, J. Appl. Phys. **39**, 3305 (1968).
- ¹⁵ P. Gas and F. d’Heurle, in *Properties of Metal Silicides*, edited by K. Maex and M. Van Rossum (INSPEC, 1995), p. 279.
- ¹⁶ Y. Hiraoka, H. Iwasawa, T. Inoue, M. Nagae, and J. Takada, J. Alloy. Compd. **377**, 127 (2004).
- ¹⁷ G. Samsonov and A. Epik, ZB. PR. Inst. Teploenerg. Akad. Nauk Ukr. R.S.R. **26**, 67 (1964).
- ¹⁸ C. J. Rosa, Metall. Trans. A **14**, 199 (1983).
- ¹⁹ B. M. Warnes and G. Simkovich, J. Less-Common Met. **106**, 241 (1985).
- ²⁰ J. Nakanishi, A. Otsuki, T. Oku, and O. Ishiwata, J. Appl. Phys. **76**, 2293 (1994).

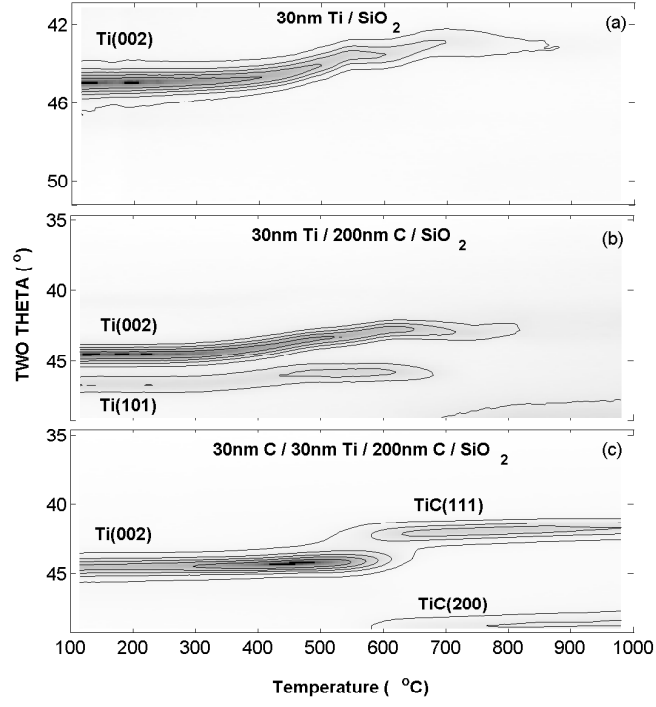


FIG. 1: *In-situ* XRD-pattern of 30nm Ti deposited (a) on SiO_2 , (b) on 200nm amorphous carbon on SiO_2 , and (c) on 200nm amorphous carbon on SiO_2 with 30nm carbon capping layer.

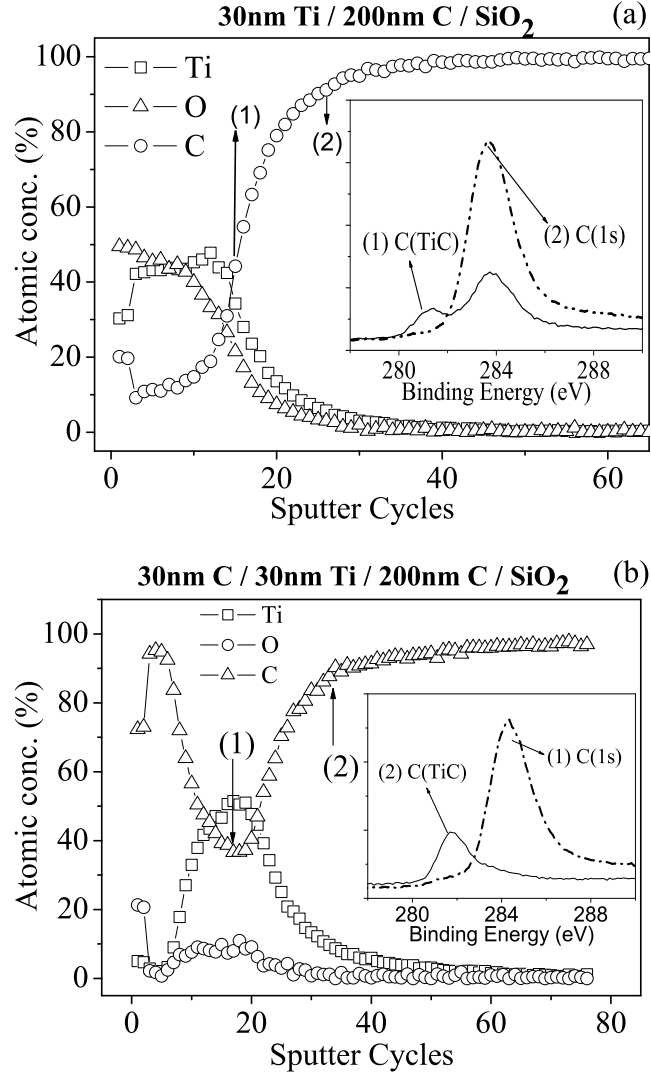


FIG. 2: XPS depth profiles of the Ti-C system: (a) without a capping layer, (b) with a 30nm C capping layer. Both samples were annealed to 1000°C at 3°C/s in a He-atmosphere. The insets show the single C-peak measured at the respective depths, marked with (1) and (2).

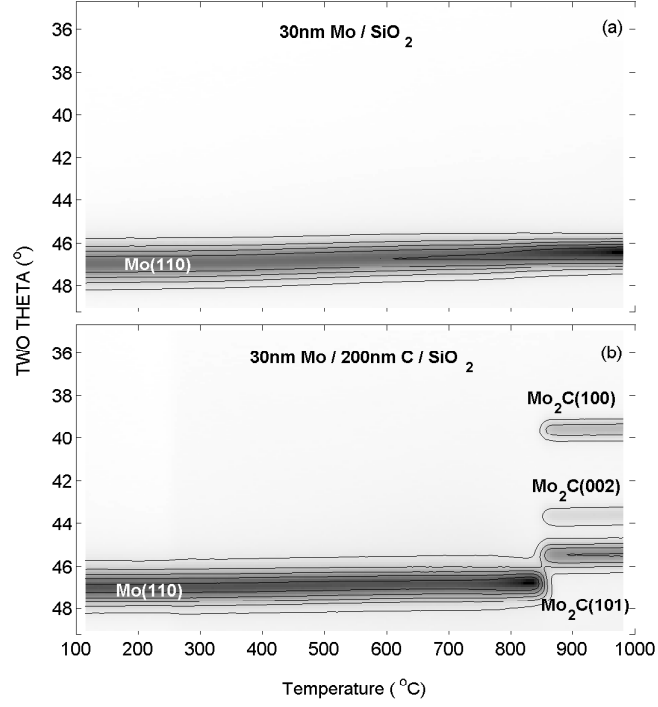


FIG. 3: *In-situ* XRD patterns of 30nm Mo deposited on SiO_2 (a), and on 200nm C on SiO_2 without a capping layer (b). Samples were annealed to $1000^\circ C$ at $3^\circ C/s$.

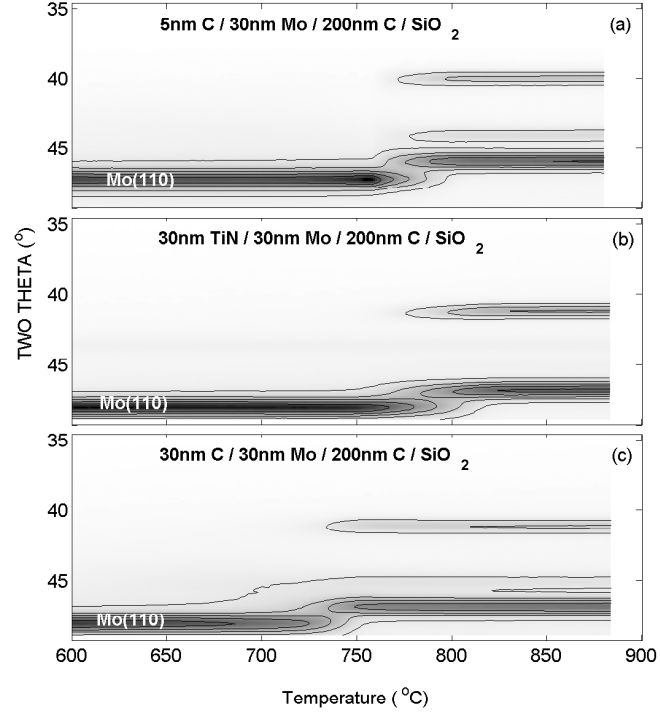


FIG. 4: *In-situ* XRD patterns of 30nm Mo deposited on 200nm C on SiO_2 with: (a) 5nm C capping layer, (b) 30nm TiN capping layer, and (c) 30nm C capping layer. The samples were annealed to 1000°C at $3^\circ\text{C}/\text{s}$.

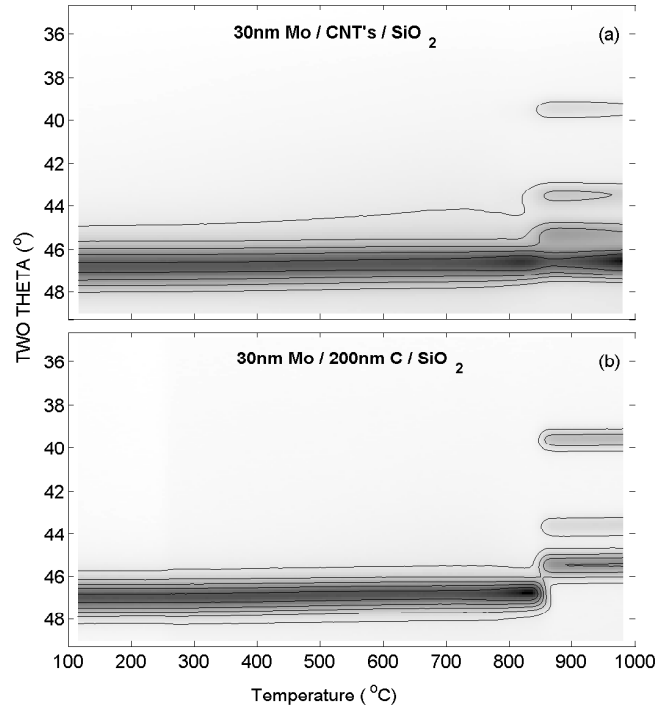


FIG. 5: 30nm of Mo deposited (a) on CNT's on SiO_2 , and (b) on 200nm C on SiO_2 (without a capping layer). The samples were annealed to $1000^\circ C$ at $3^\circ C/s$.

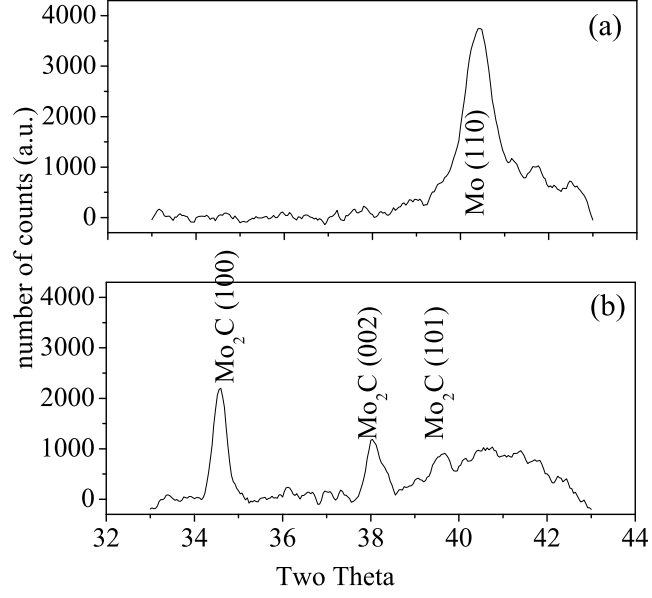


FIG. 6: Standard $\theta/2\theta$ ($CuK\alpha$) of 30nm Mo deposited on CVD diamond: (a) unreacted Mo, and (b) annealing to $1000^\circ C$ (without a capping layer) results in Mo_2C -formation.