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CFD study of droplet atomisation using a binary nozzle in fluidised bed coating

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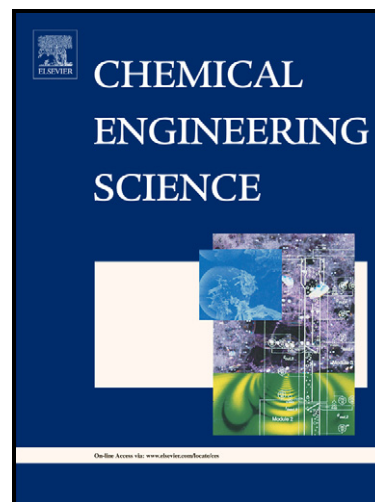
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26 behaviour with the inclusion of sprays obtained by both gas-solid-liquid multiphase modelling  
27 methods is presented.

28

29 **Keywords**

30 Fluidisation, Computational Fluid Dynamics, Multiphase flow, Powder technology, Two-fluid  
31 Atomisation, Population balance model

32

33 **1. Introduction**

34 Among a wide range of microencapsulation techniques, fluidised bed technology has been  
35 successfully used for the coating of particulate solids due to its excellent mixing capabilities  
36 and its optimal heat and mass transfer rates (Ronsse *et al.*, 2007a, b). Fluidised bed coating is  
37 an added-value technique whereby a pure active ingredient or mixture of ingredients, in solid  
38 particulate form, is encapsulated within a coating polymer. The aim of encapsulation is to  
39 control release, to protect the core ingredients, to increase the overall product quality and to  
40 increase the processing convenience. An aqueous or organic solvent-based solution containing  
41 the coating polymer is continuously sprayed by means of a pneumatic or two-fluid nozzle,  
42 which may be submerged in or positioned above the bed (Depypere *et al.*, 2009; Ronsse *et al.*,  
43 2007b). In top-spray configuration, regarded as the most appropriate method to be  
44 successfully used in the food industry due to its high versatility, relatively high batch size and  
45 relative simplicity (Depypere *et al.*, 2009), the two-fluid nozzles are usually positioned above  
46 the bed, producing sprays of an aqueous solution of the coating material with a droplet size  
47 ranging from 10 to 40  $\mu\text{m}$  in order to coat particles (Hede *et al.*, 2008; Ronsse *et al.*, 2007b).

48

49 In top-spray fluidised bed coating, the basic operating principle consists of air suspension of  
50 particles in the coating chamber, spraying of coating polymer solution as droplets with the

51 objective of increasing the probability of particle-droplet impact, spreading of droplets on the  
52 particle surface, droplet evaporation and layering or superposition of droplets on the particle  
53 surface resulting in a homogeneous coating enveloping the core particles (Teunou and  
54 Poncelet, 2002). In order to control process efficiency in fluidised bed coating using a model-  
55 based approach, it is necessary to explore each phenomenon taking place in the system. As  
56 described in previous works (Duangkhamchan *et al.*, 2010; Duangkhamchan *et al.*, 2011), the  
57 momentum transfer between the gas and solid phases was first modelled using various drag  
58 coefficient models, in order to evaluate the appropriate drag model for the description of  
59 fluidised bed behaviour (Duangkhamchan *et al.*, 2010). However, in that work, only  
60 interaction between gas and solid phases with the absence of atomisation was taken into  
61 account. Subsequently, the solids volume fraction was simulated including the effect of the  
62 release of compressed air by the two-fluid nozzle in order to provide qualitative and  
63 quantitative consistency of model simulations with experimental data (Duangkhamchan *et al.*,  
64 2011). However, the liquid phase, being the sprayed droplets, was not yet included in the  
65 latter study. Therefore, the next step – as outlined in this research article – is the addition of  
66 the liquid phase to the existing fluidised bed CFD model.

67

68 To produce sprays in fluidised bed coating processes, pneumatic or two-fluid atomisation is  
69 frequently used. In the mechanism of the two-fluid atomisation, as shown in Fig.1, a high  
70 velocity gas impacts a liquid jet issuing from a nozzle orifice creating high shear force over  
71 the liquid surface, leading to disintegration into spray droplets. The optimum frictional  
72 conditions resulting from high relative velocity between gas and liquid are generated by  
73 expanding the air to sonic or supersonic velocities before impacting the liquid (Hede *et al.*,  
74 2008). When injected from the nozzle orifice, the liquid jet starts to make contact with the  
75 mixing zone, expanding radially and squeezed into a thin circular sheet (Zeoli and Gu, 2006).

76 The term “liquid sheet” is used for both flat and cylindrical jets as common nomenclature  
77 (Hede *et al.*, 2008). For more details about two-fluid atomisation, the reader is referred to  
78 Hede *et al.* (2008), Sridhara and Raghunandan (2010) and Varga *et al.* (2003).

79

80 [Insert Figure 1 here]

81

82 Currently, to design and optimise the fluidised bed coating process, spray conditions and the  
83 operation of the two-fluid nozzle are identified as one of the most critical factors for the  
84 whole process and in practice have to be trial-and-error tested in order to control spray  
85 characteristics, including droplet size distribution, droplets trajectories and spray cone angle  
86 (Hede *et al.*, 2008; Ronsse *et al.*, 2007b). Therefore, in order to reduce time consumption and  
87 expensive cost of extensive experiments, many numerical approaches, for instance, Eulerian-  
88 Eulerian CFD, Eulerian-Lagrangian CFD, and population balance modelling, etc., have been  
89 developed as a powerful tool to comprehend or clarify the impact of different input variables  
90 on process efficiency and to research and design work (Ronsse *et al.*, 2007b).

91

92 During the last few decades, CFD has been widely adopted in many industrial uses. In spray  
93 application, various numerical methods, for instance, the volume of fluid (VOF) method and  
94 the discrete phase method (DPM), have been developed to predict basic characteristics of  
95 spraying nozzles (e.g., spray angle and droplet size distribution) and to predict droplet  
96 trajectories. In the discrete phase method (Lagrangian framework), the droplet trajectory is  
97 calculated individually using the equation of motion, whereas the volume of fluid or  
98 multifluid method (Eulerian framework) is based on continuum mechanics which treat the  
99 two phases as interpenetrating continua (Taghipour *et al.*, 2005).

100

101 For instance, for studying the two-fluid atomisation, instead of using only experimental PIV  
102 to provide an instantaneous map of the entire velocity field, Hoeg *et al.* (2008) used Eulerian  
103 CFD models to investigate the flow pattern of gas and liquid jets issuing from a two-fluid  
104 nozzle. In that work, good agreement between model-predicted and experimental data was  
105 found. Furthermore, Zeoli and Gu (2006) used the discrete phase model to simulate the  
106 critical droplet breakup during atomisation producing fine spherical metal powders. To verify  
107 their model performance, the liquid metal was initialised to large droplet diameters varying  
108 from 1 to 5 mm. They found that the model could provide quantitative assessment for the  
109 atomisation process. Pimentel *et al.* (2006) improved the capability of CFD models to capture  
110 liquid atomisation mechanisms of the two-fluid nozzle associated with the measured droplet  
111 diameters to initialise the droplet size in the discrete phase model. Even though many  
112 researchers have attempted to model droplet atomisation in various applications by means of  
113 CFD, as seen in Behjat *et al.* (2010), Fuster *et al.* (2009), Gianfrancesco *et al.* (2010), Kalata  
114 *et al.* (2009), Mezhericher *et al.* (2010), White *et al.* (2004) and Yamada *et al.* (2008), two-  
115 fluid atomisation occurring in the fluidised bed coating process still needs to be explored,  
116 considering the fact that the liquid is atomised in the presence of the fluidised solid phase.

117

118 In addition to the two approaches for modelling multiphase flow problems mentioned  
119 previously, the population balance model has been introduced to the CFD community as an  
120 alternative approach because of its reduced level of computational complexity (Aly *et al.*,  
121 2009). Moreover, the model can be easily coupled with the Eulerian-Eulerian model which  
122 eliminates the need for semi-empirical models employed in the Lagrangian framework (Aly *et*  
123 *al.*, 2009). Recently, the population balance model has been extensively used in liquid-liquid  
124 and gas-liquid systems for modelling droplets and bubbles (Aly *et al.*, 2010a, b). However,  
125 only few studies can be found in droplet atomisation problems, especially in fluidised bed

126 coating systems. The atomisation process occurring in a plain jet air blast atomiser (two-fluid  
127 nozzle) was first investigated using the combination between a population balance model and  
128 a CFD Eulerian multi-fluid model by Aly *et al.* (2009). In that work, although the model  
129 obtained good agreement with experimental data, improvement still needed to be done.  
130 Therefore, Aly *et al.* (2010a, b) developed a new mathematical model for calculating droplet  
131 breakup frequency, instead of using a constant value, based on both drag and turbulence  
132 induced fragmentation stresses. Good agreement with the experimental data was achieved.

133

134 The main objective of this work is to present a CFD model of droplet atomisation of a two-  
135 fluid nozzle in the fluidised bed coating process, and to integrate it with existing gas-solid  
136 CFD models for fluidised bed coating processes as described by Duangkhamchan *et al.* (2010,  
137 2011). Furthermore, the alternative numerical approach to describe two-fluid atomisation  
138 using population balance modelling combined with the Eulerian CFD framework is also  
139 demonstrated. Finally, the impact of process variables on spray characteristics and  
140 comparison of model-predicted distribution of voidage and liquid volume fraction obtained by  
141 two approaches are assessed.

142

143 The results presented in this paper are part of a research project aiming at modelling the  
144 complete top-spray fluidised bed coating process using CFD, with the global aim of  
145 understanding the process fundamentals and to provide the insight for optimising process  
146 control and reactor design. The modelling of the complete coating process requires several  
147 aspects to be studied in more detail and consequently, the research was split up into four parts,  
148 pertaining to the modelling of these aspects. First, a CFD model with appropriate selection of  
149 a drag model was constructed to allow the accurate prediction of gas/solid behaviour in  
150 tapered fluidised beds (Duangkhamchan *et al.*, 2010). Next, the effect of the release of

151 compressed air – to assist in the atomisation of the coating solution – on the hydrodynamic  
152 behaviour of the fluidised bed was studied (Duangkhamchan et al., 2011). The third part and  
153 also the subject of this research paper, deals with the hydrodynamic modelling of the liquid  
154 spray in the gas/solid fluidised bed. Finally, the overall CFD model will be concluded by  
155 adding the heat and mass transfer (i.e. evaporation of the binder solution in the droplets and as  
156 deposited onto the particles).

157

## 158 **2. CFD model description**

### 159 ***2.1. Discrete phase model (DPM)***

160 In addition to solving transport equations for the continuous phases (i.e. gas and solids), a  
161 discrete phase of droplets was simulated in a Lagrangian framework. The trajectories of these  
162 discrete phase entities were computed individually. The coupling between the phases and its  
163 impact on both the discrete phase trajectories and the continuous phase flow was included.

164

#### 165 *2.1.1 The Euler – Lagrangian approach*

166 In the Euler-Lagrangian approach, the gas and solid phases are treated as continuous phases  
167 by solving the time-averaged Navier-Stokes equations, while the dispersed phase (liquid  
168 phase) is solved by tracking a large number of droplets through the calculated flow field. The  
169 discrete phase can exchange momentum with the fluid phase.

170

#### 171 Continuous phase model

172 Each volume within the mesh is simultaneously solved in an Eulerian frame of reference to  
173 obtain the gas flow field with the use of general conservation equations, as summarised below  
174 and described in more detail in Duangkhamchan *et al.* (2010, 2011).

175



176 The conservation of mass of phase  $q$  ( $q =$  either gas or solid) is described as

177

$$\frac{\partial}{\partial t}(\alpha_q \rho_q) + \nabla \cdot (\alpha_q \rho_q \vec{v}_q) = 0 \quad (1)$$

178

179 where  $\alpha_q$  is the phase volume fraction,  $\rho_q$  the density and  $\vec{v}_q$  the velocity of phase  $q$ .

180

181 The following equation describes conservation of momentum for the fluid phase  $l$ :

182

$$\frac{\partial}{\partial t}(\alpha_l \rho_l \vec{v}_l) + \nabla \cdot (\alpha_l \rho_l \vec{v}_l \vec{v}_l) = -\alpha_l \nabla \cdot \vec{\tau}_l + \alpha_l \rho_l \vec{g}_l + K_{ls}(\vec{v}_l - \vec{v}_s) \quad (2)$$

183

184 In Eq. (2) is  $(\vec{v}_l - \vec{v}_s)$  the slip velocity between the phases, where the subscript  $l$  denotes the  
 185 fluid phase and  $s$  indicates the solid phase.  $K_{ls}$  denotes the drag force coefficient relevant to  
 186 the phases  $l$  and  $s$ ,  $p$  stands for the pressure and  $\vec{\tau}_l$  the deviatoric effective stress tensor of  
 187 fluid phase.

188

### 189 Discrete phase model

190 The discrete phase model was solved in a Lagrangian frame of reference to simulate the spray  
 191 pattern which is predicted by tracking the droplet trajectories by integrating the force balance  
 192 of the particle (Newton's second law) (Behjat *et al.*, 2010; Pimentel *et al.*, 2006),

193

$$\frac{d\vec{v}_p}{dt} = \vec{F}_D + \frac{\vec{g}(\rho_p - \rho_p)}{\rho_p} \quad (3)$$

194

195 where  $\vec{F}_D$  is the drag force per unit particle mass,

196

$$\vec{F}_D = \frac{3}{4} \cdot \frac{C_D \rho_g}{\rho_p d_p^2} \cdot |\vec{v}_g - \vec{v}_p| \cdot (\vec{v}_g - \vec{v}_p) \quad (4)$$

197

198 Here,  $\vec{v}_g$  is the gas phase velocity,  $\vec{v}_p$  is the droplet velocity,  $\rho_g$  is the gas density,  $\rho_p$  is the  
 199 droplet density and  $d_p$  is the droplet diameter.  $C_D$  is the drag coefficient estimated using the  
 200 correlation proposed by Morsi and Alexander (1972).

201

### 202 2.1.2 Air-blast/air-assisted atomiser model

203 Among the five atomiser models available in FLUENT, namely the plain-orifice, pressure-  
 204 swirl, flat-fan, effervescent/flashing and air-blast/air-assisted atomiser models, the latter was  
 205 found to be the suitable model for the two-fluid nozzle in this work due to similar atomisation  
 206 mechanism. The air-blast/air-assisted atomiser model predicts droplet formation in those  
 207 atomisers where an additional air stream is used to accelerate the breakup of the liquid sheet  
 208 formed by the nozzle into droplets. In this case, droplet formation is characterised by the  
 209 production of a liquid sheet, which further breaks up into ligands which finally disintegrate  
 210 into droplets as shown in Figure 1.

211

212 In order to determine the liquid sheet thickness, the effective mass flow rate,  $\dot{m}_{\text{eff}}$ , defined as

213

$$\dot{m}_{\text{eff}} = \frac{2\pi \dot{m}}{\Delta\phi} \quad (5)$$

214

215 is used. In this equation,  $\Delta\phi$  is the difference between the azimuthal stop angle and the  
 216 azimuthal start angle, which was  $2\pi$  for the nozzle type considered in this study (circular

217 liquid sheet). Here,  $\dot{m}$  is the liquid mass flow rate ( $\text{kg s}^{-1}$ ) and hence, was equal to  $\dot{m}_{\text{eff}}$ .

218 Therefore, the thickness of sheet (mm),  $d_{\text{sh}}$ , produced by the air-blast/air-assisted atomiser can

219 be approximated by relating the mass flow rate as

220

$$\dot{m}_{\text{eff}} = \pi \rho_1 v_1 d_{\text{sh}} (d_{\text{inj}} - d_{\text{sh}}) \quad (6)$$

221

222 where  $\rho_1$  is the liquid density ( $\text{kg m}^{-3}$ ),  $v_1$  is the axial velocity component of the liquid at the

223 nozzle orifice ( $\text{m s}^{-1}$ ) and  $d_{\text{inj}}$  is the diameter of the liquid orifice. In Eq. (6), the effective mass

224 flow rate is expressed as the liquid density ( $\rho_1$ ) multiplied with the liquid velocity ( $v_1$ ) and

225 with the cross-sectional area of the circular liquid sheet with diameter ( $d_{\text{inj}} - d_{\text{sh}}$ ) and sheet

226 thickness  $d_{\text{sh}}$ .

227

228 Owing to the instability of the liquid sheet, the sheet breaks up and ligaments will be formed

229 (Figure 1) whose length is given by

230

$$L_{\text{lg}} = C_{\text{sh}} \frac{v_{\text{sh}}}{\Omega} \quad (7)$$

231

232 where  $L_{\text{lg}}$  is the ligament length (mm),  $C_{\text{sh}}$  denotes a sheet constant assumed to be responsible

233 for sheet breakup,  $v_{\text{sh}}$  is the total velocity of the liquid sheet and  $\Omega$  is the maximum growth

234 rate ( $\text{s}^{-1}$ ) and is found by numerically maximising the dispersion relation based upon the

235 growth of sinuous waves on the liquid sheet (Schmidt *et al.*, 1999). For short waves, the

236 ligament diameter is assumed to be linearly proportional to the wavelength that breaks up the

237 sheet,

238

$$d_{lg} = \frac{2\pi C_{lg}}{K'} \quad (8)$$

239

240 where  $C_{lg}$  is the ligament constant, and  $K'$  is the wave number ( $m^{-1}$ ) corresponding to the  
 241 maximum growth rate,  $\Omega$ . For more details concerning the air-blast/air-assisted atomiser  
 242 model, the reader is referred to Ansys Inc. (2009b).

243

## 244 **2.2 Multi-fluid flow model combined with population balance model (MFM-PBM)**

245 In this approach, a conservation equation called the population balance equation is solved  
 246 along with the Navier-Stokes equations in order to calculate the droplet diameter (Sauter  
 247 Mean Diameter, SMD, which is commonly used to characterise droplets in spray modelling)  
 248 and subsequent transport throughout the atomisation process. The population balance equation  
 249 is a statement of continuity that describes how the statistical distribution of one or more  
 250 droplet-related variables changes with time and space (Peglow *et al.*, 2007). If the fraction of  
 251 droplets with volume  $V$  at time  $t$  is given by the number density function  $n(V, t)$ , then the  
 252 change in number of droplets with volume  $V$  as a result of larger droplets with volume  $V'$ ,  
 253 fragmenting into droplets with volume  $V$  and the subsequent fragmentation of droplets with  
 254 volume  $V$  into smaller droplets, is given in the population balance equation as (Ansys Inc.,  
 255 2009a):

256

$$\frac{\partial}{\partial t} [n(V, t)] + \nabla \cdot [\vec{v}_p n(V, t)] = \int_V g(V') \beta(V/V') n(V', t) dV' - g(V) n(V, t) \quad (9)$$

257

258 In the above equation, only droplet fragmentation or break-up was considered while droplet  
 259 coalescence was assumed to be negligible since sprays in the fluidised bed coating process are  
 260 considered to be dilute. Furthermore, droplet breakup was assumed to be binary, i.e., when a

261 droplet with volume  $V'$  breaks up, it forms two new droplets with volume  $V$  and  $(V'-V)$ . The  
 262 terms on the left hand side of Eq. (9) are the rate of change of the number density function  
 263 and its convective derivative, respectively, while the terms on the right hand side represent the  
 264 birth rate and the death rate terms resulting from droplet breakage (Ansys Inc., 2009a).

265

266 In Eq. (9)  $g(V')$  is breakage frequency, being the fraction of droplets of volume  $V'$  breaking  
 267 per unit time ( $s^{-1}$ ) and  $\beta(V/V')$  is the droplet breakage kernel and expresses the probability that  
 268 a droplet with volume  $V$  originates from the binary fragmentation of a droplet with volume  $V'$ .  
 269 Hence, the first term on the right hand side of Eq. (9),

270

$$\int_{V'} g(V') \beta(V/V') n(V', t) dV' \quad (10)$$

271

272 represents the rate of formation, or birth rate, of droplets with volume  $V$  from breakage of  
 273 droplets with volume  $V'$  ( $V < V' < \infty$ ). The second term on the right hand side of Eq. (9),

274

$$g(V) n(V, t) \quad (11)$$

275

276 expresses the rate at which droplets with volume  $V$  disappear from the system due to  
 277 fragmentation in smaller droplets, hence the term death rate.

278

279 Many methods, like for instance the Monte Carlo method (Lasheras *et al.*, 2002; Ronsse *et al.*,  
 280 2007a), the discrete method (Aly *et al.*, 2009; Lasheras *et al.*, 2002), the quadrature method of  
 281 moments (Marchisio *et al.*, 2003), and the direct quadrature method of moments (Madsen,  
 282 2006), are widely used in order to solve the population balance equations.

283

284 Due to the advantages of computing the droplet size distribution directly and the assumption  
 285 of a small number of size intervals and the size distribution, in this work the discrete method  
 286 (also known as the class method as referred to in this paper) was used to discretise the droplet  
 287 population into a finite number of size intervals,  $n$ . In the class method, assuming that  
 288 aggregation is negligible for dilute sprays, the population balance equation is written in terms  
 289 of volume fraction of droplets with size class  $i$  (Aly *et al.*, 2010; Aly *et al.*, 2009):

290

$$\frac{\partial}{\partial t}(\rho_p \alpha_i) + \nabla \cdot (\rho_p v_i \alpha_i) = \rho_p V_i (B_i - D_i) \quad (12)$$

291

292 where  $V_i$  is the size of droplets in size class  $i$ ,  $\rho_p$  is the density of the droplet phase and  $\alpha_i$  is  
 293 the volume fraction of droplet size class  $i$  defined as

294

$$\alpha_i = N_i V_i \quad i = 0, 1, \dots, n-1 \quad (13)$$

295

296 where  $N_i$  is the total number of droplets per size class, and is given by

297

$$N_i = \int_{V_i}^{V_{i+1}} n(V) dV \quad (14)$$

298

299 The droplet birth rate resulting from droplet breakage,  $B_i$ , and death rate term,  $D_i$ , in Eq.(12),  
 300 are defined as

301

$$B_i = \sum_{j=1}^n g(V_j) \beta(V_i/V_j) N_j \quad (15)$$

$$D_i = g(V_i) N_i \quad (16)$$

302

303 In this study, breakage kernel and breakage frequency in the death and birth rate terms in Eqs.  
304 (15) and (16) were similar to those used in Aly *et al.* (2009): the breakage kernel  
305 corresponded to a case where equal droplet fragments are distributed to all daughter size bins,  
306 while a constant frequency of 2000 Hz – a number corresponding to the reciprocal of the  
307 mean characteristic time scale of turbulence eddies – was chosen for the breakage frequency.

308

309 In order to solve the number density function, the population balance equation is linked to the  
310 Eulerian CFD model via a two-way coupling procedure. In this procedure, the velocity  $v_i$  is  
311 calculated in the Eulerian framework and then substituted into the population balance  
312 equation in order to compute the mean droplet sizes (SMD) which are then returned to the  
313 Eulerian solver to calculate the phase interaction such as the momentum exchange. Heat and  
314 mass transfer were not taken into account in this work.

315

### 316 **3. Materials and methods**

#### 317 ***3.1. Overview of the numerical method***

318 In this work, in addition to the discrete phase model (DPM) which has been widely used to  
319 describe droplet atomisation, the basic class method of population balance model combined  
320 with the multi-fluid phase model (MFM-PBM) was used to simulate the two-fluid  
321 atomisation. For both approaches, the 3-D geometry of a laboratory-scale Glatt GPCG-1  
322 fluidised bed (Glatt GmbH, Germany) together with the two-fluid nozzle were meshed using  
323 Gambit 2.2.30 (Ansys Inc., Canonsburg, PA). A hybrid hexahedral-tetrahedral grid,  
324 containing 473 083 elements, was exported into the solver software, Ansys Fluent v.12  
325 (Ansys Inc., Canonburg, PA). The grid is displayed in Fig. 2.

326

327 Flow turbulence was simulated using the standard  $k-\varepsilon$  model with standard wall functions,  
328 which has proven to result in good agreement with experimental data for this type of fluidised  
329 bed reactor geometry and solids material, as demonstrated in Duangkhamchan *et al.* (2010,  
330 2011). First order upwind schemes were selected for the convection terms and the relation  
331 between velocity and pressure corrections was calculated using the phase-coupled SIMPLE  
332 algorithm.

333

334 In order to model the droplet atomisation and the droplets' interaction with the gas and solid  
335 phases, the numerical setup was separated into four sections corresponding to two approaches  
336 as follows:

337

- 338 • modelling droplet atomisation in a gas phase using the discrete phase model
- 339 • modelling droplet atomisation in a gas-solid flow by means of the discrete phase  
340 model combined with the multifluid flow model
- 341 • modelling droplet atomisation in a gas phase using the population balance model  
342 combined with the multifluid flow model
- 343 • modelling droplet atomisation in a gas-solid flow using the population balance model  
344 coupled with the multifluid flow model.

345

346 [Insert Figure 2 here]

347

### 348 **3.2. Modelling the two-fluid atomisation using the DPM**

#### 349 *3.2.1. Two-phase flow model (gas-liquid DPM model)*

350 First, model parameters associated with the air-blast/air-assisted atomiser model were  
351 calibrated using published industrial data of the mean droplet diameter for the nozzle used in



352 the Glatt GPCG-1 fluidised bed unit (Model 970-S1, Düsen-Schlick GmbH). In order to  
353 compare the calculated droplet sizes to the published industrial data, spray injections (water)  
354 at atomisation air pressures of 1.0 and 3.0 bar were simulated in a stagnant-air cylinder  
355 geometry.

356

357 Multiple simulations with variation of spray sheet thicknesses ( $d_{sh}$ ) of 0.4, 0.7 and 1.0 mm,  
358 sheet constants ( $C_{sh}$ ) of 14, 15 and 16, and ligament constants ( $C_{lg}$ ) of 0.5, 0.7 and 0.9, were  
359 performed. After calibration of the air-blast atomiser model, the model was employed to study  
360 the atomisation in the geometry of the laboratory-scale fluidised-bed. Fluidisation air flow  
361 rates similar to those used in the previous works (Duangkhamchan *et al.*, 2010, 2011), namely  
362 55, 76 and 97 m<sup>3</sup>hr<sup>-1</sup>, were selected. Other process variables are presented in Table 1.

363

364 [Insert Table 1 here]

365

### 366 3.2.2. Three-phase flow model (gas-solid-liquid DPM model)

367 A CFD model including all three phases and their interactions – i.e. momentum transfer – was  
368 developed to evaluate the impact of injection parameters. In this model, the water droplets  
369 were considered to be a separate phase in addition to the gaseous and solid particle phases in  
370 order to better describe the complex process of liquid spray inside the fluidised bed reactor.  
371 Interactions between gas and solid phases were solved in the Eulerian-Eulerian framework,  
372 including the modified-Gidaspow drag coefficient as employed in Duangkhamchan *et al.*  
373 (2010), while the trajectories of injected droplets were simulated by solving the equations of  
374 motion of individual dispersed phase entities. For fluidisation, the properties of 1 kg of glass  
375 beads were used (Depypere *et al.*, 2009) as listed in Table 2, along with the boundary  
376 conditions and simulation parameters.

377

378 [Insert Table 2 here]

379

380 **3.3. Modelling two-fluid atomisation using the MFM-PBM**381 *3.3.1. Two-phase flow model (gas-liquid MFM-PBM model)*

382 The unsteady pressure-based solver in Ansys Fluent v.12 was used with the Eulerian  
383 multiphase model coupled with the discrete population balance model in order to simulate  
384 water droplet atomisation through the two-fluid nozzle. Assuming that sprays are dilute,  
385 turbulence was modelled using the dispersed k- $\epsilon$  turbulence model. To ensure solution  
386 convergence behaviour within each time step, a small time step of  $10^{-5}$  s was chosen. The  
387 velocity boundary conditions for water and atomising air are shown in Table 1. As the use of  
388 compressed air in the two-fluid nozzle produces droplets with a size ranging from 10 to 40  
389  $\mu\text{m}$  (Lefebvre, 1989), in this work, the droplet population was discretised into 7 size classes  
390 with a diameter ranging from 10 to 40  $\mu\text{m}$ . The breakage kernel was computed to represent a  
391 case where droplet fragments are distributed to all daughter size bins.

392

393 *3.3.2. Three-phase flow model (gas-solid-liquid MFM-PBM model)*

394 Instead of solving the motion of injected droplets separately in the Lagrangian discrete phase  
395 model, the gas, droplets and solid particles were treated as interpenetrating continua in the  
396 Eulerian framework. The gas phase was considered to be the primary phase, whereas the  
397 droplets and solid particle phases were the secondary phases. To track the droplet diameter in  
398 the Eulerian solver, the number density function was solved using the class population  
399 balance method (see Section 2.2).

400

401 **3.4. Experimental spray visualisation**

402 *3.4.1. Measurement set-up*

403 Spray visualisation experiments were performed in a transparent, polycarbonate reactor with  
404 similar dimensions to the Glatt GPCG-1 fluidised bed coating reactor. The tapered reactor had  
405 a bottom diameter of 0.15 m, a top diameter of 0.30 m and a total height of 0.56 m. The  
406 reactor shell material consisted of 5 mm thick polycarbonate. Fluidisation air was provided by  
407 a 2.2 kW high pressure centrifugal fan (Ventomatic CHT160-2T-3, Belgium) equipped with  
408 electronic frequency control (Figure 3). The volumetric air flow rate was measured between  
409 the fan and the reactor inlet by means of a 0.1 m diameter rotating vane flow meter (Airflow  
410 Developments, VMD20, UK).

411

412 The air distributor used in the fluidised bed reactor consisted of a Robusta 172×36 wpi (wires  
413 per inch) wire mesh (Spörl KG, Germany). The pneumatic nozzle (Schlick Model 970-S1,  
414 Germany), normally used in the Glatt GPCG-1, was installed in the lower nozzle port of the  
415 transparent reactor. The spraying liquid was water with an added fluorescent dye, being  
416 sodium fluorescein salt (Sigma-Aldrich) and was transported to the pneumatic nozzle by  
417 means of a peristaltic pump (Watson-Marlow, 505 Du/RL, US).

418

419 The spray cone produced by the pneumatic nozzle was visualized by means of UV  
420 illumination by directing a 400 W UV spotlight 0.5 m above the reactor outlet. Illumination  
421 through the open reactor outlet proved to be most efficient, as the polycarbonate reactor shell  
422 material had some UV absorbing capacity. The illuminated spray was recorded by a digital  
423 camera (Olympus i-Speed 1) at 60 fps (800 × 600 pixel size), and stored in an uncompressed  
424 video format (AVI).

425

426 [Insert Figure 3 here]

427

428 *3.4.2. Data processing*

429 To visualise the spray cone, post-processing of the captured image data was necessary (i.e.  
 430 contrast enhancement). The uncompressed frames captured by the video camera can be  
 431 considered to be matrices holding the pixel intensity values (between 0 and 1),  $A_i$ , with  
 432 subscript  $i$  indicating the frame number (i.e. 60 per recorded second of video) and having  
 433 dimensions of 800 by 600. As the recording is triggered at the moment of activating the  
 434 peristaltic pump, an initial number of frames,  $b$ , is recorded without spray, i.e., before the  
 435 development of the actual spray cone. Of these  $b$  frames, an average,  $A'_b$ , is calculated which  
 436 serves as a reference blank frame which will be subtracted from the actual spraying frames,

437

$$A'_b = \frac{1}{b} \sum_{i=1}^b A_i \quad (17)$$

438

439 The actual frames to be used for visualisation were taken after the steady-state spraying cone  
 440 had developed, which are denoted by frames  $A_i$  with  $m \leq i \leq n$ . The contrast enhancement of  
 441 each of the frames during steady-state spraying consisted of two steps. First, the reference  
 442 blank frame was subtracted from the spraying frame  $A_i$  resulting in  $A''_i$ , and second, each  
 443 frame matrix was multiplied with a scalar so the pixel with the highest intensity value in  $A_i$   
 444 reached unity (i.e. maximum intensity) – with the new resulting matrix denoted as  $E_i$ :

445

$$m \leq i \leq n \rightarrow \begin{cases} A''_i = A_i - A'_b \\ E_i = x \cdot A''_i \Leftrightarrow \max(E_i) = 1 \end{cases} \quad (18)$$

446

447 Finally, the average matrix of all  $E_i$  with  $m \leq i \leq n$  was taken and used as the contrast-  
 448 enhanced spray visualisation image,

449

$$E' = \frac{1}{n-m} \sum_m^n E_i \quad (19)$$

450

## 451 **4. Results and Discussion**

### 452 **4.1. Air-blast/air-assisted atomiser model calibration**

453 Several simulations with different spray injection setups were performed in order to calibrate  
454 the air-blast atomiser model as outlined in Section 3.2.1. The measured mean volume droplet  
455 diameters at 1.0 and 3.0 bar atomisation pressure (data supplied by Düsen-Schlick GmbH)  
456 were compared with simulated mean droplet sizes while the ligament constant, the sheet  
457 constant and sheet thickness were varied.

458

459 Figure 4 shows the simulated versus experimental mean droplet sizes at 3.0 bar atomisation  
460 pressure and using different values for the ligament constant,  $C_{lg}$ . It can be observed that the  
461 spray injection characterised by using a ligament constant,  $C_{lg}$  of 0.9, gave the best agreement  
462 with the experimental mean droplet diameter. With respect to variation of the sheet constant  
463 ( $C_{sh}$ ) and the sheet thickness ( $d_{sh}$ ), no significant differences in model-predicted droplet  
464 diameter distribution were observed. Similar trends were seen from simulations with 1.0 bar  
465 atomisation pressure and consequently, the following calibrated air-blast atomiser model  
466 parameters were used in subsequent simulations: Ligament constant,  $C_{lg} = 0.9$  and standard  
467 values for sheet constant,  $C_{sh} = 15$  and sheet thickness,  $d_{sh} = 0.7$  mm.

468

469 [Insert Figure 4 here]

470

### 471 **4.2. Two-phase (gas-liquid) model of droplet atomisation in the fluidised bed coater**

#### 472 **4.2.1. Discrete phase model (gas-liquid DPM)**

473 Using calibrated atomiser model parameters, the spray was modelled using DPM inside the  
474 geometry of a fluidised bed coater, as detailed in Table 2 and Figure 2. The effects of  
475 fluidisation air flow rate, atomisation air pressure and liquid feed rate – as outlined in Table 1  
476 – on spray characteristics were simulated.

477

478 Figure 5 demonstrates the contour plots of droplet mass fraction at different atomisation air  
479 pressures. It can be seen that the spray pattern did not change with increasing atomisation air  
480 pressure. When considering the impact of fluidisation air flow rate, as detailed in Figure 6,  
481 higher flow rates were seen to reduce the diameter of the spray cone and to lower the droplet  
482 mass fraction in the reactor. This can be explained by the fact that, at higher fluidisation air  
483 flow, droplets are easier lifted out of the reactor. However, it is important to stress that the  
484 DPM model in its current state did not include droplet evaporation (no energy equation).  
485 Consequently, the length of model-predicted droplet trajectories is likely to be overestimated  
486 compared to the actual process, where droplets are subjected to spray drying. Finally, the  
487 effect of liquid feed rate on the droplet mass fraction distribution is shown in Figure 7. At  
488 higher liquid feed rates, more droplets can be produced. Due to the higher amount of spray  
489 issuing from the nozzle in Figure 7c, the spray cone shape can obviously not be characterised.

490

491 [Insert Figure 5 here]

492 [Insert Figure 6 here]

493 [Insert Figure 7 here]

494

495 *4.2.2. Multi-fluid flow model combined with population balance model (gas-liquid MFM-*  
496 *PBM)*

497 Figure 8 shows the distribution of the Sauter mean droplet diameter (SMD) and droplet  
498 volume fraction in the geometry of the fluidised bed coater, as predicted by the MFM-PBM  
499 gas-liquid model under reference scenario conditions as outlined in Table 1. From this figure,  
500 it can be observed that the mean droplet sizes decreases continuously as the droplets depart  
501 from the liquid orifice of the nozzle until they reach the air distributor.

502

503 [Insert Figure 8 here]

504

505 Comparison of the droplet volume fraction as shown in Figure 8b with the droplet mass  
506 fraction from the DPM model gives a moderately good agreement (same order of magnitude  
507 if mass fraction is converted to volume fraction). The exception is a more narrow spray cone  
508 that was predicted by the MFM-PBM model. Furthermore, as the droplets were treated as a  
509 continuum in the MFM-PBM model, the droplet phase was seen to deflect from the air  
510 distributor at the bottom of the reactor (Figure 8b). This effect is not visible as such in the  
511 DPM results (Figure 5b), because in the DPM algorithm, droplet tracking was ended when a  
512 droplet impacted on the boundaries of the reactor geometry. Given these results, the MFM-  
513 PBM can be opted for as an alternative approach to model the atomisation of the two-fluid  
514 nozzle, considering advantages including minimum level of computational complexity, ease  
515 of coupling with the Eulerian-Eulerian CFD model and eliminating the need for semi-  
516 empirical models employed in the DPM model (Aly *et al.*, 2009).

517

518 Simulations with varying liquid feed rate, atomisation air pressure and fluidisation air flow  
519 rate were also performed. However, significant differences in droplet size distribution were  
520 not predicted by the MFM-PBM model with respect to process condition dependency. This  
521 observation is not consistent with both numerical and experimental studies reported in

522 literature (Hede *et al.*, 2008; Lal *et al.*, 2010; Lebas *et al.*, 2009; Liao and Lucas, 2009; Liu *et*  
523 *al.*, 2006; Sridhara and Raghunandan, 2010). The possible explanation of this inconsistency is  
524 the assumption of a constant breakage kernel and breakage frequency in this study. In reality,  
525 breakage depends on droplet properties including size, as well as on local turbulences in the  
526 flow field. Consequently, improvement in model accuracy of the MFM-PBM model could be  
527 achieved if the breakage frequency and kernel are made dependant on (i.e. 'sensitised to') the  
528 droplet Weber and local Reynolds numbers (Aly *et al.*, 2009; Aly *et al.*, 2010a, b).

529

### 530 **4.3. Spray pattern validation (gas-liquid two-phase flow)**

531 Spray visualisation images combined with UV illumination as a function of atomisation air  
532 pressure are presented in Fig. 9. The effect of fluidisation air flow rate, varied within the  
533 interval outlined in Table 1, showed no observable difference in spray pattern. Regarding  
534 atomisation air pressure, higher air pressure not only resulted in a narrower spraying cone, but  
535 also resulted in a slightly asymmetric spraying cone, as shown in Fig. 9d. The asymmetry is  
536 likely due to the higher volume of expanding atomisation air (for  $P_{at} = 3$  bar) in combination  
537 with the asymmetric construction of the reactor resulting from the one-sided nozzle support  
538 (not shown in Fig. 9).

539

540 [Insert Figure 9 here]

541

542 When the experimental results are compared with the model-predicted spray patterns, a  
543 qualitative agreement was obtained with the DPM or MFM-PBM predicted results. However,  
544 DPM gave better agreement compared to MFM-PBM including a wider spraying cone and  
545 much less deflection of the spray at the base of the reactor. As already stated in Section 4.2.2.,  
546 as opposed to the MFM-PBM, the DPM can effectively account for droplet-wall collisions



547 which corresponds with the actual (experimental) process, where most of the droplets were  
548 seen to adhere onto the wire mesh air distributor at the base of the reactor.

549

#### 550 **4.4. Three-phase (gas-solid-liquid) model of droplet atomisation in the fluidised bed coater**

##### 551 *4.4.1. Discrete phase model (gas-solid-liquid DPM)*

552 Figure 10a shows the contour plot of the time-averaged steady-state voidage, taken over a  
553 simulated time period of 10s. The initial 5s of the simulated process were discarded to avoid  
554 the start-up fluidisation behaviour. As can be observed in Fig. 10a, in the central part of the  
555 reactor, the region under the nozzle is occupied by the hollow atomisation cone. In this zone,  
556 the solid particles have to be lifted by the fluidising air against the counterforce of the  
557 atomisation air resulting in a voidage (Duangkhamchan *et al.*, 2011), while a denser zone can  
558 be noticed in a radial area between the nozzle atomisation air cone and the reactor walls. It  
559 could be explained that particles move predominantly upwards in the centre to the above bed  
560 region (about 12 cm high above the air distributor), then move radially towards the walls and  
561 downwards along the walls. This particle flow behaviour was confirmed by experimental  
562 results obtained by Positron Emission Particle Tracking (PEPT) (Depypere *et al.*, 2009).

563

564 [Insert Figure 10 here]

565

566 Fig. 10b confirms that the contacting between the droplet and fluidising particles occurs at the  
567 central part of the vessel. Fig. 10b demonstrates the droplets tracked at 15s. The calculated  
568 droplet tracks revealed that droplets moved downwards along with the atomisation air cone  
569 until facing the counter-current fluidising solid particles. Considering the absence of  
570 phenomena including droplet evaporation and droplet/solids adhesion, the DPM algorithm  
571 continues to track the droplets until they exit the reactor at the top or impact one of the reactor

572 geometry boundaries. In reality, the majority of the droplets adhere onto the fluidised  
573 particles, contributing to the layered growth of the coating wall around the individual core  
574 particles. Also, in an actual fluidised bed coating process, the majority – typically  $\geq 70\%$   
575 (Ronsse *et al.*, 2008) – of water in the coating solution is evaporated after the droplets have  
576 impacted the surface of the fluidised core particles. Only a minority of the water is evaporated  
577 during droplet travel between the nozzle and the impacting particle surface. Consequently,  
578 there will be a minimal impact of the droplet size reduction as a result of droplet drying on the  
579 droplet dynamics (i.e. altered drag force, reduced droplet mass) and the resulting droplet  
580 trajectories.

581

582 *4.4.2. Multi-fluid flow model combined with population balance model (gas-solid-liquid*  
583 *MFM-PBM)*

584 Comparison between the gas-solid-liquid phase DPM and MFM-PBM model-predicted  
585 results is shown in Figure 10a-d. As can be seen in Fig. 10a and 10c, the model-predicted  
586 time-averaged steady-state voidage profiles of both models have a strong agreement. Also,  
587 when considering the model-predicted distribution of droplets within the fluidised bed (Figs.  
588 10b and 10d), and specifically the penetration depth of the droplets in the bed, the MFM-PBM  
589 predicted results are consistent with those from the DPM.

590

## 591 **5. Conclusions**

592 As a powerful numerical tool for solving fluid flow problems, CFD was used to model the  
593 important aero- and hydrodynamic aspects of a fluidised bed coater, including the gas, liquid,  
594 and solid phases using two approaches: Eulerian-Lagrangian and combined Eulerian-  
595 Eulerian/population balance model. In the discrete phase model (DPM), the calibrated air-  
596 blast/air-assisted atomiser model was used as and the effects of process variables on spray

597 flow and its mass distribution were studied. It was, in the gas-liquid DPM model, that the  
598 spray cone and liquid mass fraction change with the variation of fluidisation air flow rate,  
599 atomisation air pressure and liquid feed rate.

600

601 The population balance model combined with the Eulerian-Eulerian CFD model was  
602 employed as an alternative approach to describe the two-fluid atomisation and the impact of  
603 process variables. The gas, droplets and solid particles were modelled by treating all phases as  
604 interpenetrating continua in the Eulerian framework, while the class population balance model  
605 was used to track the droplet diameter. The simulated results showed that even though this  
606 approach could be opted for instead of the DPM model to capture the two-fluid atomisation  
607 and interaction between phases, improvement of the population balance model by for instance  
608 including more accurate breakage kernels (i.e. depending on Weber and Reynolds numbers)  
609 has to be carried out. When evolving from a gas-liquid to a gas-solid-liquid CFD model,  
610 consistency between the DPM and the population balance model in the Eulerian framework  
611 was shown to improve. Consequently the MFM-PBM approach was considered to be a viable  
612 alternative in the CFD modelling of gas-solid-liquid systems, including fluidised bed coaters.

613

614 Finally, given the absence of thermodynamics in the presented model, effects such as droplet  
615 evaporation could not be captured with the CFD model in its current state. Future work will  
616 comprise the addition of the energy equation for the description of heat transfer phenomena,  
617 as well as droplet evaporation and to tie together all phenomena occurring in the gas-liquid-  
618 solid multiphase system into a single comprehensive CFD model, suitable for accurately  
619 describing fluidised bed coating processes with the aim of improving process understanding  
620 and optimising the coating process in terms of process conditions and reactor design.

621

622 **Nomenclature**

623	$A, A', A'', E$	Image intensity matrices
624	$B$	Birth rate, $s^{-1}$
625	$C$	Constant
626	$D$	Death rate, $s^{-1}$
627	$d$	diameter, m
628	$F$	Force per unit mass, $N\ kg^{-1}$
629	$\bar{g}$	Acceleration due to gravity, $m\ s^{-2}$
630	$g()$	Breakage frequency, $s^{-1}$
631	$K$	Interphase momentum exchange coefficient, $kg\ m^{-3}\ s^{-1}$
632	$K'$	Wave number corresponding to the maximum growth rate, $m^{-1}$
633	$L$	Length, m
634	$\dot{m}$	Mass flow rate, $kg\ s^{-1}$
635	$n()$	Number density function
636	$p$	pressure, Pa
637	$Re$	Reynolds number, dimensionless
638	$t$	Time, s
639	$v$	velocity, $m\ s^{-1}$
640	$V, V'$	Volume, $m^3$
641	$\dot{V}$	Volumetric flow rate, $m^3\ s^{-1}$
642	$x$	Scalar
643		
644	<b>Greek symbols</b>	
645	$\alpha$	Volume fraction, dimensionless
646	$\beta()$	Probability density function, kernel function

647	$\rho$	Density, $\text{kg m}^{-3}$
648	$\tau$	Deviatoric stress tensor, $\text{kg m}^{-1}\text{s}^{-2}$
649	$\mu$	Dynamic viscosity of gas, $\text{kg m}^{-1}\text{s}^{-1}$
650	$\phi$	Angle, rad
651	$\Omega$	Maximum grow rate, $\text{s}^{-1}$

652

653 ***Subscripts***

654	at	atomisation
655	q	solid or gas phase
656	b	initial
657	D	drag
658	eff	effective
659	f	fluidisation
660	g	gas phase
661	i	class or integer
662	inj	injector exit
663	l	fluid phase
664	s	solid phase
665	lg	ligament
666	p	particle or droplet
667	sh	sheet

668

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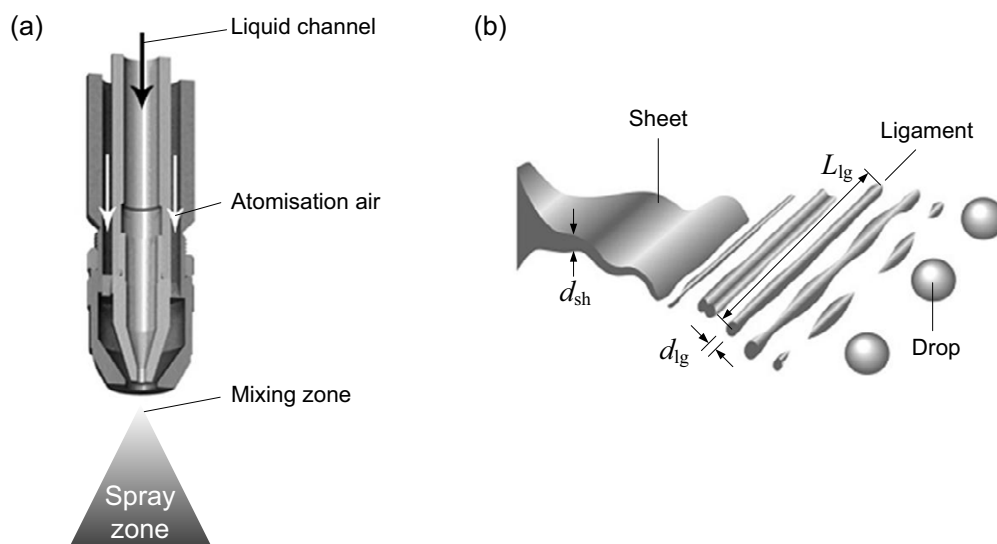
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812 **List of Tables and Figures**

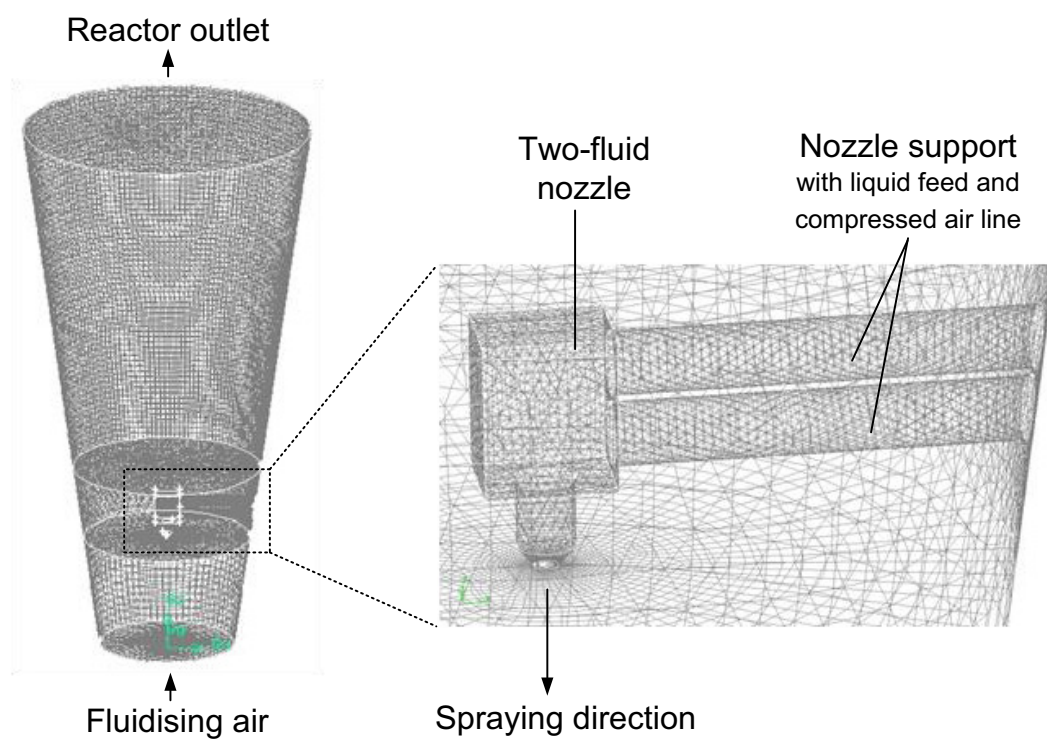
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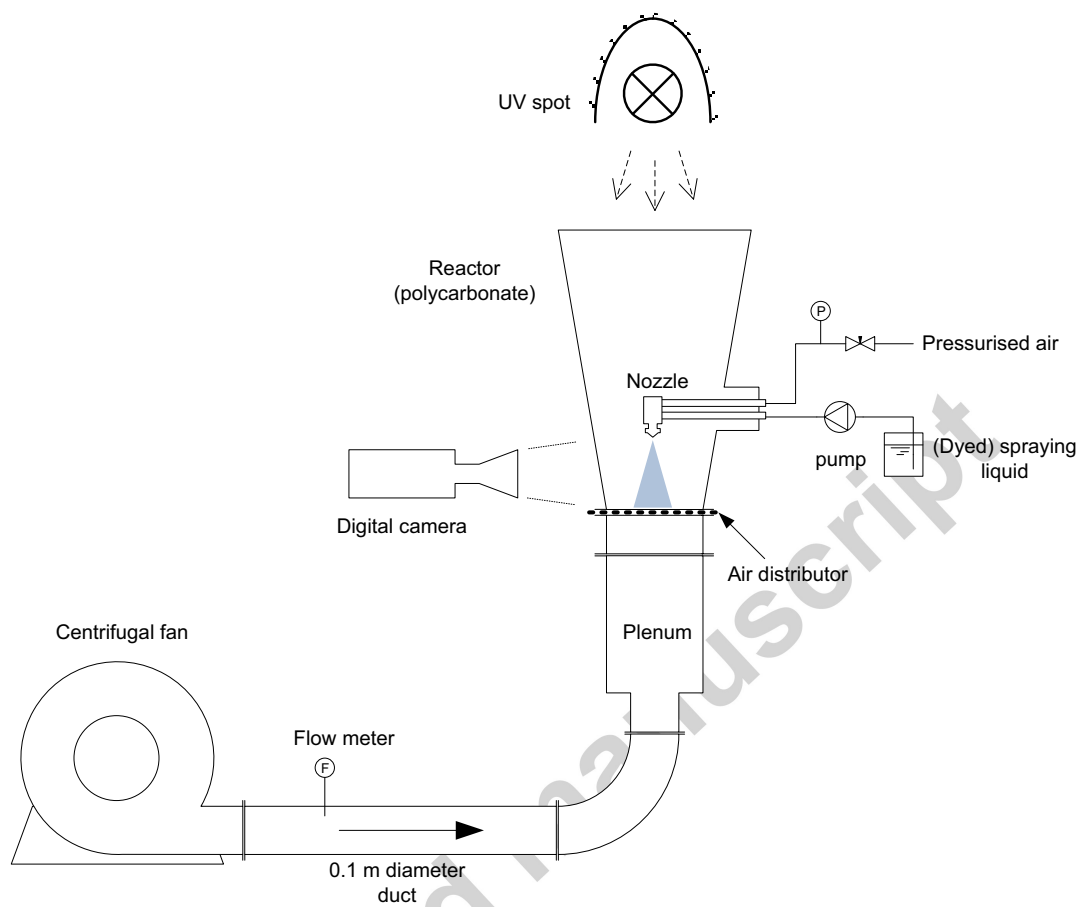
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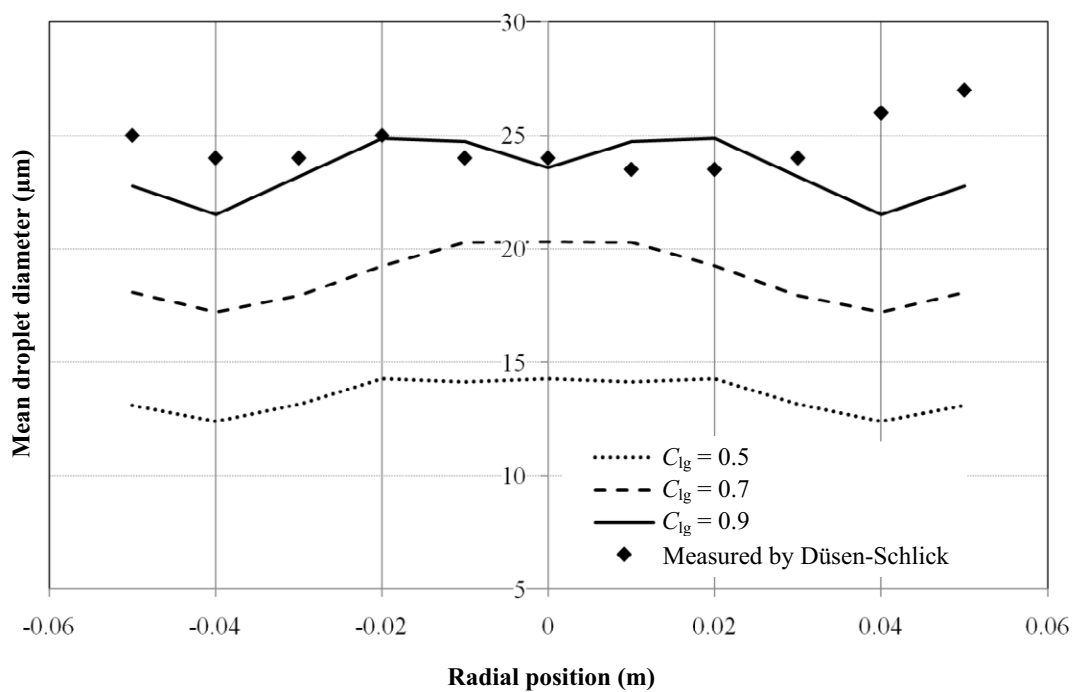
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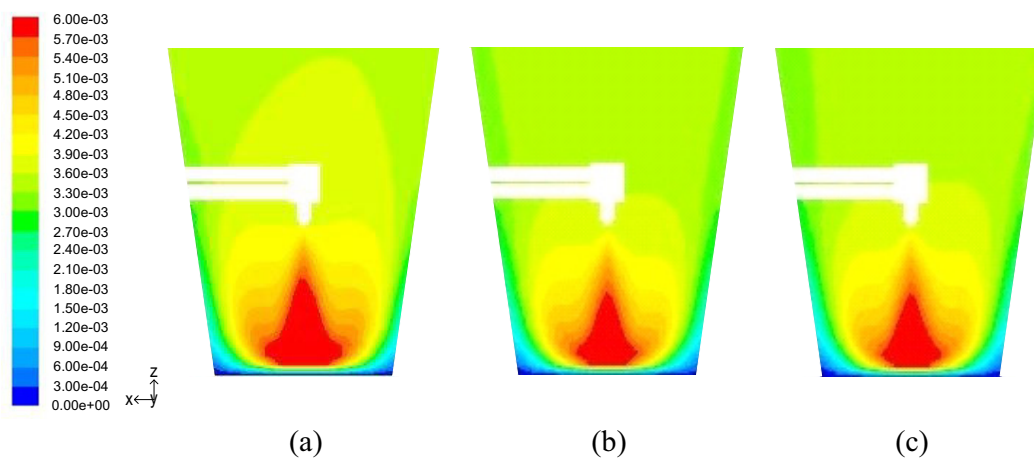
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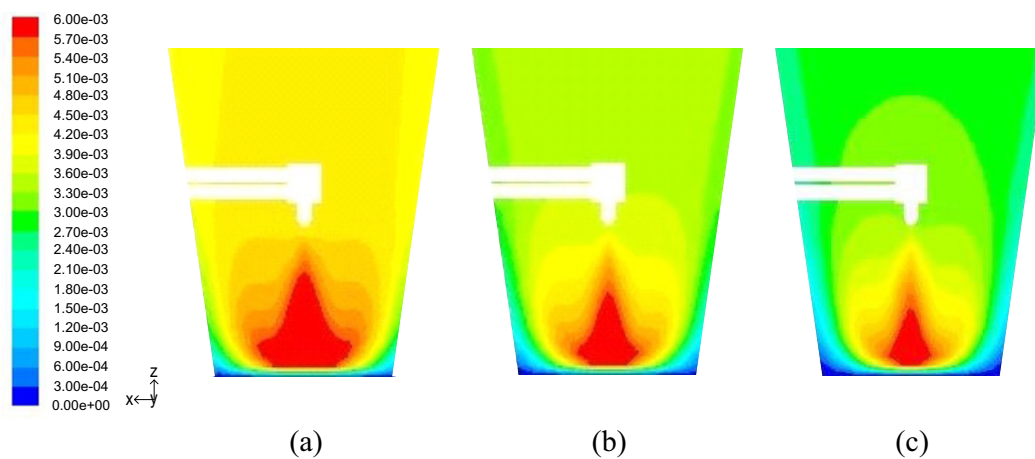
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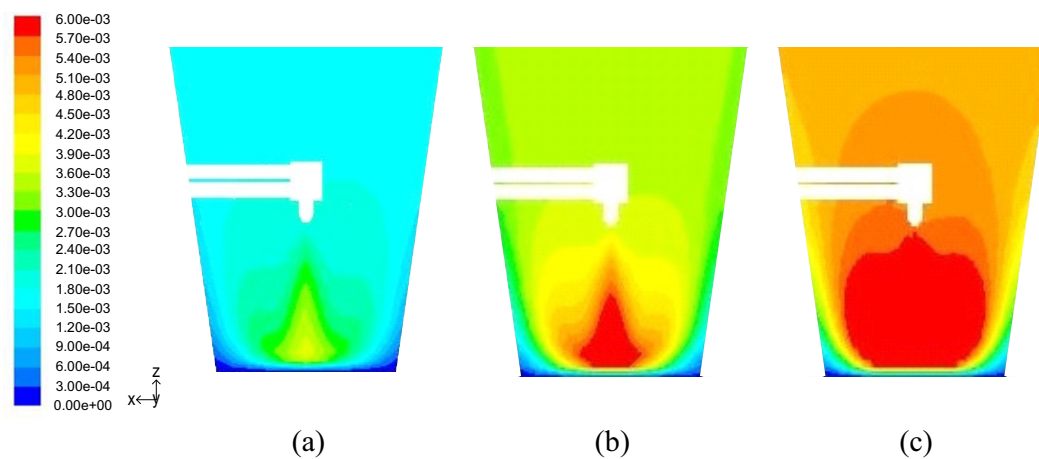
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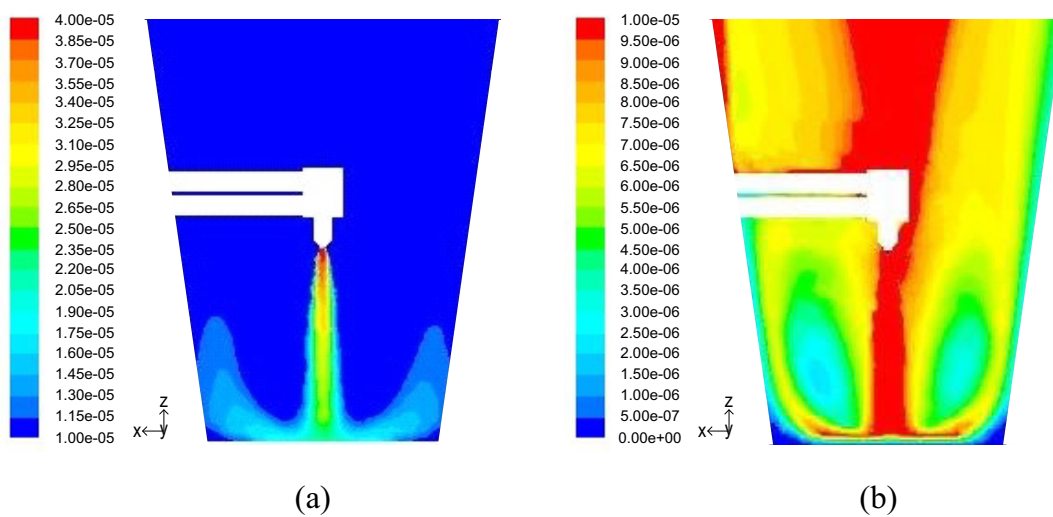
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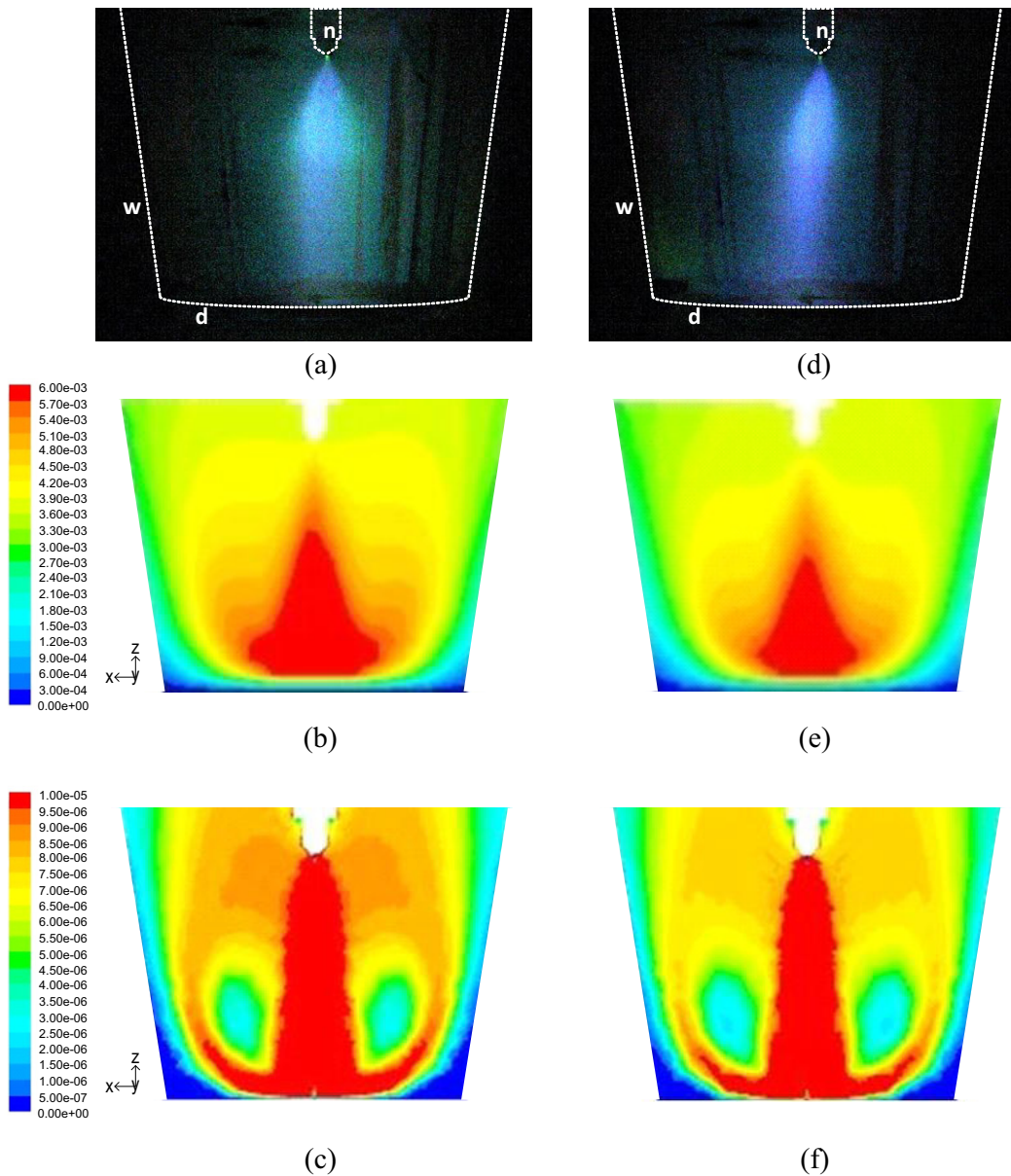
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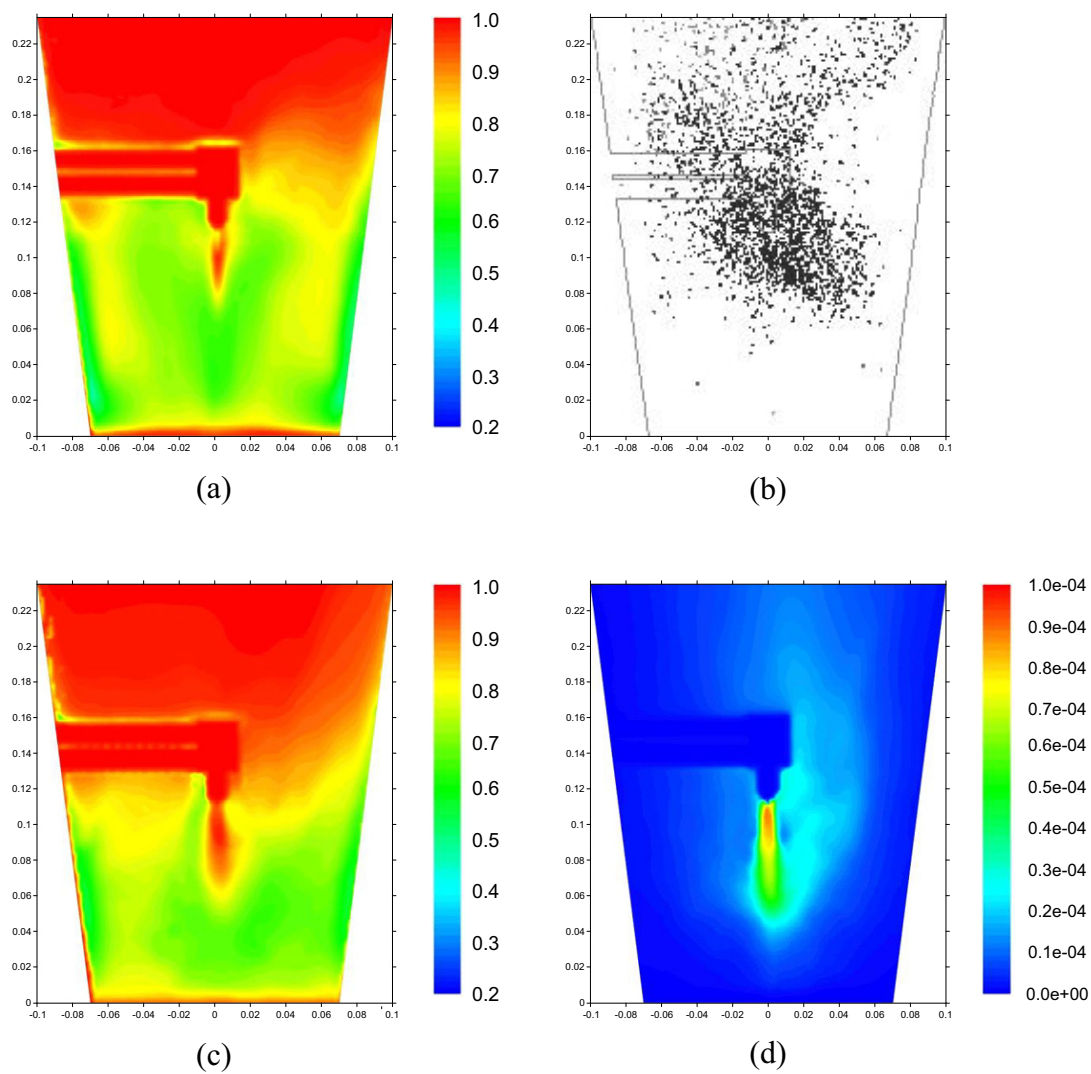
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**Figure 10.** Comparison between (a) contour of gas-solid-liquid DPM model-predicted time-averaged steady-state voidage, (b) gas-solid-liquid DPM model-predicted droplet tracks at  $t = 15$ s, (c) contour of gas-solid-liquid MFM-PBM model-predicted time-averaged steady-state voidage and (d) contour of gas-solid-liquid MFM-PBM model-predicted time-averaged steady-state liquid volume fraction.



**Table 1.** Process variables used in this study (<sup>a</sup> indicates reference scenario).

Process variables	Value
Fluidisation air flow rate, $\dot{V}_f$ (m <sup>3</sup> hr <sup>-1</sup> )	55, 76 <sup>a</sup> , 97
Atomisation air pressure, $P_{at}$ (bar)	1.0, 2.0 <sup>a</sup> , 3.0
Liquid feed rate, $\dot{M}_l$ ( $\times 10^{-4}$ kg s <sup>-1</sup> )	0.5, 1.0 <sup>a</sup> , 1.5

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**Table 2.** Boundary conditions and simulation parameters.

Descriptor	Value
Primary phase (continuous)	Gas
Secondary phase (continuous)	Glass beads
Discrete phase	Water droplets
Solids particle size, $d_s$ ( $\mu\text{m}$ )	196.54
Solids density, $\rho_s$ ( $\text{kg m}^{-3}$ )	2467
Solids loading, $M_s$ (kg)	1
Gas phase density, $\rho_g$ ( $\text{kg m}^{-3}$ )	1.225
Liquid phase density, $\rho_p$ ( $\text{kg m}^{-3}$ )	998
Reactor bottom diameter (m)	0.15
Reactor top diameter (m)	0.30
Reactor height (m)	0.56

## Highlights:

- Multiphase computational fluid dynamics model was built for fluidised bed coating
- Gas-solid fluidisation modeled in the Eulerian framework
- Two-fluid atomisation was described by discrete phase and population balance models
- Gas-liquid model-predicted spray pattern experimentally verified using UV illumination
- Population balance modelling proved viable alternative to discrete models

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