



# Géotechnique

---

## **Modelling the infiltration behavior of foam into saturated sand considering capillary resistance for EPB shield tunneling**

GEOT-2021-146 | Paper

Submitted by: Dongzhu Zheng, Adam Bezuijen, Markus Thewes

Keywords: EPB, PHYSICAL MODELLING, FOAM INFILTRATION, FOAM-SOIL INTERACTION, CAPILLARY RESISTANCE

PDF auto-generated using **ReView**  
from





# 1 **Modelling the infiltration behavior of foam into saturated sand considering** 2 **capillary resistance for EPB shield tunneling**

3 Dongzhu Zheng <sup>a,\*</sup>, Adam Bezuijen <sup>a,b</sup>, Markus Thewes <sup>c</sup>

4 <sup>a</sup> *Department of Civil Engineering, Ghent University, Ghent, Belgium*

5 <sup>b</sup> *Deltares, Delft, The Netherlands*

6 <sup>c</sup> *Institute for Tunnelling and Construction Management, Ruhr-University Bochum, Bochum, Germany*

7 <sup>\*</sup> *Corresponding author address: Technologiepark 68, Ghent, 9052, Belgium, Tel. +32 9 2645724, Fax. +32*  
8 *9 2645849, e-mail: dongzhu.zheng@ugent.be*

## 9 **Abstract**

10 With reference to the Earth Pressure Balance (EPB) shield tunneling, the pressure infiltration  
11 of foam into saturated sand is numerically investigated based on experimental analysis. The  
12 model accounts for foam spurt during foam infiltration that was experimentally discovered in a  
13 previous paper. The maximum penetration depth by foam bubbles is estimated with a simplified  
14 micro stability model based on the minimum pressure difference over an individual foam  
15 bubble across the pore throats. Although the micro stability model underestimates the maximum  
16 penetration depth in one of the sands used, it predicts well with the results obtained for the other  
17 two sands. Possible reason for the underestimation in one of the sands is discussed. Further  
18 results from numerical simulation are in accordance with the measured discharge behavior  
19 during foam spurt. The general agreement suggests that the model could explain the foam flow  
20 behavior and can be used to describe foam spurt during foam infiltration that can be expected  
21 in EPB shield tunneling.

22 **Key words:** EPB, Physical modelling, Foam infiltration, Foam-soil interaction, Capillary  
23 resistance

## 24 **Introduction**

25 The wide application of the EPB shield tunnel-boring machines (TBM) has led to a number of  
26 studies in related subjects (Bezuijen & Dias, 2017; Budach & Thewes, 2015; Bezuijen, 2012;  
27 Thewes & Budach, 2012; Peila et al., 2009; Merritt & Mair, 2006; Anagnostou & Kovári, 1996).  
28 Owing to the development of the soil conditioning technology, the application range of the EPB

29 shield TBM has been extended from clayey soils to coarse-grained soils (Thewes, 2007).  
30 Conditioning agents like foam, bentonite and polymer are often used to condition the excavated  
31 soil into a low permeable, homogenous plastic paste in order to correctly apply the support  
32 pressure onto the tunnel face.

33 Foam is the most important additive with multiple benefits including the temporary increase of  
34 porosity and compressibility as well as decreasing the shear stress and the permeability of the  
35 soils (Thewes & Budach, 2012; Psomas, 2001; Bezuijen & Schaminée, 1999). Most research  
36 focused on the bulk properties of the conditioned soils such as fluidity (Peila et al., 2009), shear  
37 behavior (Psomas, 2001), compressibility (Bezuijen, 2013) and permeability (Borio et al., 2010)  
38 and the effect of conditioned soils on the mechanical parts of the machinery such as tool wear  
39 (Wei et al., 2019) and rotating torque (Merritt & Mair, 2006). While only limited research can  
40 be found regarding the interaction between foam and soils (Xu et al., 2021; Wu et al., 2020;  
41 Wang et al., 2020; Galli, 2016).

42 Bezuijen (2002) proposed a groundwater flow model to predict the excess pore pressure in front  
43 of the tunnel face. He compared the field measurement data and found the difference in the  
44 maximum pore pressure measured in the soil in front of a slurry shield and an EPB shield is  
45 only small. This suggests that the groundwater flow model can be used in both slurry shield and  
46 EPB shield. The course of the pore water pressure in the soil for a slurry shield was investigated  
47 by Broere and van Tol (2000, 2001) and Bezuijen et al. (2001, 2016). Bezuijen et al. (2001,  
48 2016) assume the formation of an internal and external filter cake. Slurry infiltration has been  
49 extensively studied (Xu, 2018; Zizka et al., 2018; Talmon et al., 2013; Min et al., 2013).  
50 Although foam infiltration has been studied experimentally by Xu et al. (2021), Galli (2016),  
51 Quebaud et al. (1998) and Maidl (1995), the mechanism of foam infiltration remains unclear.

52 Bezuijen and Dias (2017) developed a model to describe the pressure dissipation in the mixing  
53 chamber during standstill of an EPB shield. By fitting the field measurement data to the  
54 theoretical model, the permeability of the soils in front of the tunnel face was estimated. It was  
55 found that the permeability of the soils predicted by the model is 2 to 3 times smaller than that  
56 from field measurements. Indicating that there could be a foam infiltration layer with a high  
57 flow resistance that results in the underestimation of the soil permeability by the model of  
58 Bezuijen and Dias (2017). The model was employed by Yu et al. (2020) and reasonable

59 agreement was obtained with field measurements during both drilling and standstill.  
60 Galli et al. (2021) conducted a series of experiments on foam penetration into cohesionless soils  
61 with the aim to assess the residual water content during drilling of an EPB shield TBM. The  
62 penetration process was described qualitatively by regression analysis using the power law  
63 model within the testing duration of 3 minutes. Xu et al. (2021) compared the infiltration  
64 characteristics between foam and slurry infiltration tests and found that a foam spurt process is  
65 also present compared with mud spurt in slurry infiltration. Zheng et al. (2021) found through  
66 experiments that the foam spurt results in an internal low permeable layer that creates a large  
67 pressure drop over the infiltrated sand. Although there are still fine bubbles migrating further  
68 into the sand after foam spurt, most of the bubbles are entrapped among the sand pores near the  
69 sand surface and the dominating flow will be the drainage behavior of foam.

70 The use of foam in EPB shield tunneling is still mainly based on experience over nearly 50  
71 years of development and has not become a standard process regarding its injection strategy.  
72 One possible reason may be the lack in the understanding of the foam-soil interaction, of which  
73 foam flow in saturated sand plays an important role. With the aim to understand how the foam  
74 infiltration could influence the flow behavior during standstill as well as to what extent the low  
75 permeable layer with high flow resistance could be formed, a series of foam infiltration  
76 experiments was conducted (Zheng et al., 2021). This paper focuses on the model interpretation  
77 of the foam infiltration behavior. The model is based on an infiltration model for slurry  
78 infiltration (Bezuijen et al., 2016). In addition, a finite permeability of the water through the  
79 foam in the sand is assumed. An individual bubble is considered when it penetrates the sand  
80 and the limit equilibrium state is described in a simplified micro stability model. A calculation  
81 model is developed based on a micro stability model to predict the maximum penetration depth  
82 during ‘foam spurt’ (explained in the next section). The calculation model predicts the flow  
83 behavior during foam infiltration, which accounts for the bubble flow as well as the flow of the  
84 foaming liquid responsible for the drainage of foam. Results are compared with the  
85 experimental data and further application is discussed.

### 86 **Mechanisms during foam infiltration**

87 Analogous to the two processes of mud spurt and filter cake formation for slurry infiltration  
88 (Talmon et al., 2013), a comparable foam spurt process was observed by Xu et al. (2020) and it

89 is adopted in this study. It was found by Zheng et al. (2021) that there is a maximum penetration  
90 depth by foam bubbles during foam spurt. After foam spurt, the migration of foam bubbles  
91 almost stops, and the flow is mainly controlled by the foam infiltrated sand which have the  
92 lowest permeability of the whole system. Further infiltration process was found to be the main  
93 process after foam spurt. The further infiltration process is somewhat similar to the deep bed  
94 filtration during slurry infiltration where bentonite particles are retained in the pores that the  
95 permeability of the sand decreases slowly (Yin et al., 2021; Xu & Bezuijen, 2019). The  
96 mechanisms may be different because during further infiltration, fine bubbles will not remain  
97 in the pores due to adhesion but are entrapped in the pores due to its larger volume compared  
98 with the pore throats. During further infiltration, only a small proportion of the foam bubbles at  
99 the foam front could be able to travel further and the main flow will be the water flow through  
100 the foam infiltrated sand. Besides the bubble infiltration, there is always a liquid flow due to  
101 the drainage behavior of the foam. Consequently, there is an invisible liquid front ahead of the  
102 foam front (see Figure 1).

### 103 **Model derivation**

104 In this paper, the flow behaviors during foam spurt as well as the companion liquid flow through  
105 the foam in the sand are modelled. After foam spurt, the fine bubbles migrating during further  
106 infiltration are not included in the model because it is unable to determine the amount of these  
107 fine bubbles. The model is calibrated using the results of experiments, therefore the  
108 experimental setup is described.

109 As shown in Figure 2,  $x$  is the penetration depth by the foam bubbles and  $x'$  the distance between  
110 the foam bubble front and the invisible liquid front.  $\phi_l$  represents the piezometric head in the  
111 foam above the sand surface and  $\phi_0$  the piezometric head at the foam front. The piezometric  
112 head at the bottom of the sand column is taken to be 0.

113 In this set-up, a small cylinder (with diameter  $D_2$  and length  $L_{s2}$ ) is added beneath the large  
114 cylinder (with diameter  $D_1$  and length  $L_{s1}$ ) to create an extra flow resistance which makes the  
115 equivalent length ( $L_s$ ) of the sand column equal to 5 m. The hydraulic gradient in the sand  
116 column will be  $\Delta\phi/L_s = 1$  when the applied pressure is 50 kPa during the test. This is comparable  
117 to the hydraulic gradient predicted by the groundwater flow model described in Bezuijen (2002)  
118 with a shield diameter of 10 m and an extra pore water pressure of 50 kPa. Detailed explanation

119 can be found in papers by Xu (2018) and Zheng et al. (2021).  
 120 Figure 3 shows the experimental results of the infiltration depth against square root of time.  
 121 The infiltration depth is calculated assuming there is no liquid drainage from the foam and  
 122 therefore this depth is larger than the actual penetration depth by foam bubbles. It shows that  
 123 after the initial fast penetration of the foam bubbles, there is a slower discharge period which is  
 124 approximately linear against square root of time. It was found that the permeability of the foam-  
 125 infiltrated sand ( $x$  in Figure 2) is the lowest in the whole system. There is no impermeable layer  
 126 formed on top of the sand, but the linear part is related to the layer of low permeability formed  
 127 during foam spurt. After foam spurt, there is limited foam bubble infiltration and the dominant  
 128 flow will be the water flow from the foam. The linear line against square root of time is the  
 129 result of the increase of the foam infiltration layer and therefore can be simulated. There are  
 130 some different discharge patterns after the linear distribution period for some unknown reasons  
 131 as can be seen from Figure 3. This paper will only focus on the simulation of foam spurt which  
 132 comprises the initial fast discharge and the linear discharge against square root of time shown  
 133 in Figure 3.

134 As shown for slurry infiltration (Bezuijen et al., 2016), the penetration of slurry into the sand  
 135 does not follow Darcy's law, because slurry is a fluid with yield stress and a certain pressure  
 136 drop is necessary to push the slurry further into the sand. The experimental results from Xu  
 137 (2018) and Zheng et al. (2021) suggest that there is a maximum penetration depth ( $L$ ) when  
 138 foam is pressurized into saturated sand, indicating a yield stress in the foam flow. Assume foam  
 139 to be a Bingham fluid, the flow in porous media can be expressed as follows (Bezuijen et al.,  
 140 2016):

$$141 \quad v_f = \frac{k_f}{n} \left( \frac{d\phi}{dx} - i_t \right) \quad (1)$$

142 With  $v_f$  the penetration velocity of the foam bubbles,  $k_f$  the permeability of the sand for foam,  $n$   
 143 the porosity of the sand and  $i_t$  the threshold hydraulic gradient.

144 The threshold hydraulic gradient ( $i_t$ ) can be calculated when the maximum penetration depth  
 145 ( $L$ ) is known. Specifically, because there is a water flow after the foam has reached its maximum  
 146 penetration depth, a certain pressure difference over the foam front and the outlet at the  
 147 maximum penetration depth should be present. The threshold hydraulic gradient ( $i_t$ ) can be

148 estimated:

$$149 \quad i_t = \frac{\phi_1 - \phi_{0,end}}{L} \quad (2)$$

150 With  $\phi_{0,end}$  the piezometric head of the foam front at maximum penetration depth.

151 Considering there is always a drainage behavior during foam spurt, an extra water flow through  
152 the sandy foam should be accounted for in the total discharge.

$$153 \quad v_{fw} = k_{fw} \frac{\phi_1 - \phi_0}{nx} \quad (3)$$

154 With  $v_{fw}$  the penetration velocity of water with respect to the penetration velocity of the foam  
155 bubbles and  $k_{fw}$  the permeability of the sandy foam.

156 The pore water velocity in the sand ahead of the foam front can be calculated:

$$157 \quad v = k_s \frac{\phi_0}{n(L_s - x)} \quad (4)$$

158 With  $v$  the pore water velocity in the sand ahead of the foam front and  $k_s$  the permeability of the  
159 sand.

160 The total velocity of the water is the addition of the foam velocity and the water velocity with  
161 respect to the foam velocity:

$$162 \quad v = v_f + v_{fw} \quad (5)$$

163 At the maximum penetration depth, the bubble penetration stops ( $v_f = 0$ ), and there will only be  
164 a water flow. It comes to the following relation:

$$165 \quad k_{fw} \frac{\phi_1 - \phi_{0,end}}{L} = k_s \frac{\phi_{0,end}}{L_s - L} \quad (6)$$

166 The Equations 1~4 can be solved numerically using an explicit scheme starting from  $x = 0$  at  
167  $t=0$ .

168 The maximum penetration depth ( $L$ ) can be obtained by analyzing the micro stability of an  
169 individual foam bubble across the sand pores. Rossen and Gauglitz (1990) developed a model  
170 to describe the minimum pressure gradient of foam flow through porous media. But the pressure  
171 difference in their model is between gas and liquid which is more applicable in the field of  
172 enhanced oil recovery. While here a model to describe the difference in pore water pressure is



173 needed. Therefore, a simplified micro stability model is proposed to get  $L$ .  
174 In the experiments, one fine sand (Sand 1) and two medium sands (Sand 2 and 3) (Figure 4)  
175 were tested to represent different sand stratum. The foam expansion ratio (FER) of the foam  
176 used in the experiments were 10, 15 and 20. Grain (bubble) size distribution curves of the three  
177 types of sand (foam) are shown in Figure 4. It should be mentioned that the bubble size is  
178 measured at atmospheric pressure and calculated at an absolute pressure of 1.5 bar according  
179 to Boyle's Law. The volume of an individual bubble will increase as it travels further due to the  
180 decrease in its surrounding pore water pressure. The compressibility in this case will be  
181 discussed in the following part.

182 The filter rule derived by Terzaghi and regulated in USACE (2000) is a common guidance for  
183 filter design in dams and dikes. A filter material is placed downstream to prevent particle loss  
184 of the base material. A major criterion for the filter material is  $D_{15}/d_{85} \leq 4$  ( $D_{15}$  for the filter  
185 material and is the diameter which 15% of the filter material's mass content smaller than,  $d_{85}$   
186 for the base material). Assume the foam being a base material with the same gradation and the  
187 sand column being a filter material, the potential of particle loss from the base material into the  
188 filter material is first checked with the filter rule. A regular base material is different to foam  
189 because sand particles cannot deform like foam bubbles. If it fulfills the filter rule, there would  
190 hardly be any transportation of foam bubbles into the pore space of the sand column without  
191 changing shapes. The related values are summarized in Table 1. It shows that the combinations  
192 of different base and filter materials fulfil the filter rule, suggesting that no bubble penetration  
193 could take place without changing its shape.

194 Figure 5 shows a sketch of shape changes in a bubble penetration over an individual pore throat.  
195 The changes in shape require a pressure difference to initiate the deformation which is related  
196 to the geometry of the pore throat and the surface tension of the bubble. Here the mechanism is  
197 first discussed in general and the determination of relevant parameters will be introduced  
198 subsequently.

199 As illustrated in Figure 5, the bubble changes from position  $A$  to  $B$  and  $C$ .  $A'$ ,  $B'$  and  $C'$  represent  
200 the bottom part of the bubble at corresponding moments, respectively. The pore water pressure  
201 on top of the bubble is  $P_{w1}$  and at the bottom  $P_{w2}$ . The difference in pore water pressure on  
202 either side of the pore throat ( $P_{w1}-P_{w2}$ ) pushes the bubble downwards that the bottom part bulges

203 through the pore throat.

204 Because the surface tension resists the forward movement during the change from  $A'$  to  $B'$ , the  
 205 capillary pressure will be the largest at the narrowest position  $B'$  where the bottom part of the  
 206 bubble has the largest curvature. The movement will carry on from B to C when the pressure  
 207 difference ( $P_{w1}-P_{w2}$ ) can overcome this largest capillary pressure. Because position  $B-B'$  is very  
 208 crucial to a bubble penetration, it will be analyzed based on its equilibrium state.

209 By the definition of surface tension ( $\gamma$ ), the following relations can be obtained:

$$210 \quad P_a - P_{w1} = \frac{4\gamma}{D} \quad (7)$$

$$211 \quad P_a - P_{w2} = \frac{4\gamma}{d} \quad (8)$$

212 With  $D$  the local diameter on top of the bubble and  $d$  the throat diameter.

213 Combining Equation 7 and 8 yields:

$$214 \quad P_{w1} - P_{w2} = 4\gamma\left(\frac{1}{d} - \frac{1}{D}\right) \quad (9)$$

215 Equation (9) represents the minimum difference in pore water pressure for a bubble to travel  
 216 through an individual pore throat. Above this pressure difference, the bubble will continue to  
 217 penetrate further. At the maximum penetration depth ( $L$ ), the pressure difference will no longer  
 218 be able to push the bubble forward and the penetration will stop.

219 There can be two different conditions regarding the sizes between the bubble and the pore body  
 220 with the values listed in Table 1, as shown in Figure 6, these two conditions mainly influence  
 221 the local diameter on top of the bubble ( $D$ ).

222 1) The bubble size is equal to or smaller than the pore body (Fig. 4a) but bigger than the throat  
 223 diameter (sand 3-FER 10). In this case, a bubble can fully pass through an individual throat  
 224 before it starts the next penetration. The pressure drop over one foam bubble is equal to the  
 225 pressure difference over one pore throat.  $D$  in this case will be taken as the average diameter of  
 226 the bubbles. The diameter of the bubble is influenced by the pore water pressures at different  
 227 vertical locations. Assume  $i$  (starting from  $i=1$ ) to be the number of bubble penetrations along  
 228 the flow direction. The pressure drop in one bubble penetration is:

$$229 \quad P_{wi} - P_{w(i+1)} = 4\gamma\left(\frac{1}{d} - \frac{1}{D_i}\right) \quad (10)$$

230 With  $i$  the number of bubble penetrations across the pore throat and  $D_i$  the diameter at that point  
231 which can be calculated according to Boyle's law:

$$232 \quad D_{(i+1)} = \sqrt[3]{\frac{P + P_{wi}}{P + P_{w(i+1)}}} D_i \quad (11)$$

233 With  $P$  the atmospheric pressure.

234 The calculation scheme stops when  $P_{w(i+1)} \leq 0$ . Assuming the pore space to be equal to the  
235 diameter of the sand grains, the maximum penetration depth can be calculated by:

$$236 \quad L = MD_s \quad (12)$$

237 With  $M$  the maximum value of  $i$  and  $D_s$  the average diameter of the sand grains.

238 The calculation scheme described by Equations 10 to 12 can be carried out starting from  $i = 1$ ,  
239  $D_1 = D$  and  $P_{w1} = 50$  kPa.

240 2) The bubble size is bigger than the pore body (Fig.4b). An individual bubble needs to take up  
241 more than one pore throat at a time that the bubble flows as 'bubble trains' during penetration  
242 (Rossen, 1990). In this case, a volume factor ( $\alpha$ ) is introduced to represent the number of pore  
243 throats occupied by one bubble. The volume of bubble can be calculated when the pore water  
244 pressure is known. A simplification in this model takes the local diameter on top of the bubble  
245 to be equal to the pore space which is equal to the diameter of the sand grains, the pressure drop  
246 in one bubble penetration is:

$$247 \quad P_{wi} - P_{w(i+1)} = 4\gamma\left(\frac{1}{d} - \frac{1}{D}\right) \quad (13)$$

248 The number of bubbles along the infiltration path can be obtained by:

$$249 \quad N = \frac{\Delta P}{P_{wi} - P_{w(i+1)}} \quad (14)$$

250 With  $N$  the number of bubbles along the maximum penetration depth and  $\Delta P$  the total excess  
251 pore pressure.

252 Since the volume of bubble is inversely proportional to the absolute pressure, the volume factor  
253 for bubbles at different places along the penetration path can be obtained and the maximum  
254 penetration depth is:

$$L = D \sum_{i=1}^N \alpha \frac{P + \Delta P}{P + [\Delta P - (i-1) \frac{\Delta P}{N}]} \quad (15)$$

256 With  $\alpha$  the volume factor that is determined by the volume ratio between the bubble and the  
257 pore space.

258 Here  $d_{10}$  ( $d_{10,s}$  for sand and  $d_{10,f}$  for foam) is taken to be the characteristic values that are  
259 summarized in Table 2.

## 260 **Results and discussion**

### 261 *Maximum penetration depth*

262 The surface tension ( $\gamma$ ) of the foaming liquid was measured with a single capillary tube with an  
263 inner diameter of 0.15 mm and outer diameter of 0.25 mm and a length of 300 mm (CM  
264 Scientific, CV1525 Borosilicate Glass Round Capillaries) at room temperature of 20 °C. The  
265 measurement was also conducted with water as a reference. Measurement results are listed in  
266 Table 3.

267 The measurement result on water is very close to standard value of 0.07275 N/m (Vargaftik et  
268 al., 1983), indicating the reliability of the capillary rise method.

269 The throat diameter of the sand was determined by the measured water retention curve. Water  
270 retention curve was measured using the hanging water column test in which suction pressures  
271 were varied by adjusting the height of the water column in a stepwise manner (see Lins, 2009).

272 The testing method has the advantage of its ability to obtain very small suction pressure among  
273 a small pressure range since sand usually presents a relatively small range of air entry values  
274 (Lins, 2009). During the measurement, air will displace the water in the pore space as long as  
275 the pressure difference could overcome the capillary pressure. As indicated in Figure 5, the  
276 capillary pressure is the largest at the narrowest point with the throat diameter  $d$ . Therefore, the  
277 measured water retention curve has a relation with the throat diameter. The measurement data  
278 can be fitted to the van Genuchten model (1980):

$$279 \quad \Theta = \left[ \frac{1}{1 + (ah)^n} \right]^m \quad (16)$$

280 With  $\Theta$  the normalized volumetric water content,  $a$ ,  $n$  and  $m=1-1/n$  curve fitting soil parameters  
281 and  $h$  the pressure head.

282 The range of air entry values for each sand can be obtained statistically:

$$283 \quad P - P_w = \frac{4\gamma_w}{d} \quad (17)$$

284 with  $P$  the atmospheric pressure,  $P_w$  the suction pressure (negative) and  $\gamma_w$  the surface tension  
285 of water.

286 Take Sand 1 as an example, the air entry value is between 3.2~12 kPa and the calculated throat  
287 diameter is between 0.024~0.091 mm. The pore space for each sand is considered to be the  
288 same to its characteristic grain size. With the maximum and minimum throat diameter, a  
289 distribution curve regarding the throat diameter can be developed with different air entry values  
290 and their corresponding volumetric water contents. The results are shown in Figure 7.

291 With the distribution curve,  $d_{10,p}$  is taken to be the characteristic throat diameter and the values  
292 for each sand are listed in Table 2. The Terzaghi model assumes the throat diameter ( $d_{10,p}$ ) to be  
293 approximately one quarter of the characteristic grain size ( $d_{10,s}$ ). While the ratio of  $d_{10,p} / d_{10,s}$   
294 can be found in Table 2 and it is 0.29, 0.42 and 0.27 for Sand 1, Sand 2 and Sand 3, respectively.  
295 It can be seen that only for Sand 2, the ratio is apparently different from the Terzaghi value of  
296 0.25, which could be induced by its relative density of 70%. For Sand 1 and Sand 3, it is quite  
297 close (0.29 and 0.27 for Sand 1 and Sand 3, respectively) at the relative density of 80%.

298 Table 2 shows that the micro stability model predicts well compared with the experimental  
299 results for Sand 1 and Sand 2. The model underpredicts the maximum infiltration depth for  
300 Sand 3, probably because finer bubbles are smaller than the larger pore throats in Sand 3 and  
301 thus are simply transported. Although Table 1 shows that there is no bubble penetration without  
302 changing its shape with the filter rule, Figure 4 and Figure 7 show that a certain proportion of  
303 foam bubbles are smaller than the measured largest pore throats in Sand 3. For foam FER of  
304 10, this proportion is about 70% and about 60 % and 40 % for foam FER of 15 and 20,  
305 respectively. As a consequence, the multiple number of bubbles in a single pore space likely  
306 result in the bubbles being more or less connected, which results in a smaller pressure drop over  
307 an individual pore throat than the model prediction.

### 308 *Model simulation*

309 The permeability of the sandy foam for water ( $k_{fw}$ ) can be determined according to Darcy's Law:

$$310 \quad k_{fw} = \frac{\Delta L Q}{A_c \Delta \phi} \quad (18)$$

311 With  $\Delta L$  the distance between two adjacent pressure transducers and  $\Delta \phi$  the corresponding  
 312 difference in piezometric head. The distance ( $\Delta L$ ) between  $k_1$  and  $k_2$  is only 2 cm because the  
 313 first 2 cm is pure foam. For  $k_{2-3}$  and  $k_{3-4}$  this distance is also 2 cm while for  $k_{4-5}$  it is 4 cm.

314 At the maximum penetration depth, the bubbles will reach an equilibrium state which is  
 315 described in the micro stability model. At the equilibrium state, the measured pore pressures by  
 316 adjacent transducers can be used to calculate the permeability of the sandy foam, because there  
 317 is only a water flow.

318 Figure 8 shows the calculated permeabilities with Equation 18. The infiltration depth on the  
 319 horizontal axis was calculated assuming there was no liquid drainage from the foam and  
 320 therefore is larger than the actual penetration depth by the bubbles. It shows the calculated  
 321 values exhibit decreases in permeabilities by several orders of magnitude after the foam bubbles  
 322 penetrated the corresponding sand layers. For  $k_{1-2}$  this decrease happens at the beginning of the  
 323 test because the foam bubbles immediately penetrate the sand pores when the test starts. The  
 324 values of  $k_{2-3}$  kept constant for some infiltration depth, after which they also presented a sharp  
 325 decrease. A similar situation applies to  $k_{3-4}$ , but only after the foam front has penetrated the sand  
 326 layer sandwiched by  $k_3$  and  $k_4$ . Some fluctuations of  $k_{4-5}$  were observed at the later part of the  
 327 tests, which is attributed to the small water discharge at the later stage when the accuracy of the  
 328 pressure transducers and the discharge rate started to influence the results. It should be  
 329 mentioned that the theoretical infiltration depth at which  $k_{2-3}$  and  $k_{3-4}$  start to decrease for the  
 330 test with Sand 1 FER 10 should be 0.02 m and 0.04 m, respectively. While these two values are  
 331 0.04 m and 0.065 m as shown in Figure 8, larger than the theoretical values. It suggests that the  
 332 drainage from foam has resulted in a larger infiltration depth when calculated with the  
 333 discharged water. Similar condition applies to other tests shown in Figure 8.

334 Before the foam bubbles penetrate the sand pores, the calculated values by Equation 18 should  
 335 be constant and they should reflect the permeability of the sand sample. Table 4 shows a good  
 336 compliance when comparing the sand parameters and the calculated constant values from  
 337 Figure 8.

338 Basic requirements should be applied when employing the permeability of the sandy foam ( $k_{fw}$ )

339 calculated by Equation 18 in the numerical model:

340 1) The foam bubbles stop penetrating further. This criterion can be roughly met when the  
341 maximum penetration depth described in the micro stability model is reached. Although some  
342 local unstable state can be expected during the further infiltration process, most part of the foam  
343 bubbles in the sand are stranded and the calculated values by Equation 18 will be regarded as  
344 the permeability of the sandy foam ( $k_{fw}$ ).

345 2) The sand layer sandwiched by adjacent pressure transducers is fully penetrated by the foam  
346 bubbles. This is to ensure that the seepage length through sandy foam is equal to the distance  
347 between the adjacent transducers when criterion 1 is satisfied and can be distinguished from  
348 Figure 8. For instance, when  $k_{4-5}$  starts decreasing, then the sand layer between  $k_3$  and  $k_4$  is  
349 fully infiltrated by the bubbles because now the bubbles have reached beyond  $k_4$ .

350 Following the above discussions,  $k_{fw}$  was chosen to be around  $2 \times 10^{-8}$  m/s for the test with Sand  
351 1 FER 10 according to Figure 8. The values for other cases are listed in Table 4.

352 Another parameter that needs to be obtained for the numerical model is the permeability of the  
353 sand for foam ( $k_f$ ). Since it is unrealistic to obtain through measurement,  $k_f$  was determined by  
354 fitting the measurement data into the simulation results.

355 The fitting procedure starts from  $k_f$  and then both  $L$  (from the micro stability model) and  $k_{fw}$   
356 (from Figure 8) can be slightly adjusted if model output can better fit into the experimental  
357 results. The resulting  $L$  can be found in Table 2 and the permeabilities in Table 4. It should be  
358 mentioned that the changes in parameters are manual to get better outputs compared with  
359 experimental results.

360 Results from both experiments and model simulation are plotted in Figure 9. It turns out that  
361 the simulation yields good agreements with  $L$  of 0.09 m with tests of Sand 3. For Sand 1 and  
362 Sand 2, the maximum penetration depth ( $L$ ) fits well with the experimental results except in the  
363 test with Sand 2 FER 10, where the value of  $L$  is 0.06 m, larger than 0.026 m found in the micro  
364 stability model. A possible reason for the underestimation is that 40% of the foam bubbles in  
365 foam of FER 10 are smaller than the largest pore throats in Sand 2. The smaller bubbles can  
366 simply be transported during the test, which results in a larger penetration depth than the micro  
367 stability model prediction. While in general, the good agreements between experiments and  
368 model simulation indicate the right physics of the model. It suggests that the model can be used

369 to describe the infiltration behavior during foam spurt which can be expected during EPB shield  
370 tunneling.

371 Figure 10 shows the values of the permeabilities used in the model simulation. With a certain  
372 sand type, its permeability to foam is decreasing with an increasing FER. And the permeability  
373 of the sandy foam formed during pressure infiltration decreases with an increasing FER.  
374 Suggesting that a foam with larger foam bubbles and less water content (bigger FER) results in  
375 a lower permeability when infiltrating into the sand. With a coarser sand, its permeability to a  
376 certain foam will be larger, which is logical because the throat diameter will be larger in a  
377 coarser sand, and it will be easier for both foam bubbles and water to pass through. An  
378 interesting point worth further investigation is that all permeabilities decrease by a factor of  
379 approximately 2 when FER increases from 10 to 20. A possible reason is that the smaller liquid  
380 content in the foam of FER 20 results in the film of the bubbles being thinner, further resulting  
381 in a smaller permeability.

382 Furthermore, the model can be used to calculate the foam infiltration in the tunnel front and the  
383 water flow from the mixing chamber during standstill. To do that Equation 3 has to be changed  
384 (Bezuijen et al., 2016):

$$385 \quad v = k_s \frac{\phi_0}{nR} \quad (19)$$

386 With  $R$  the radius of the tunnel.

### 387 **Conclusions**

388 This paper presents the model interpretation of the mechanisms during foam infiltration, based  
389 on the analysis of the experiments in Zheng et al. (2021). A calculation model is established to  
390 describe the flow behavior during foam infiltration into saturated sand tailored to EPB shield  
391 tunneling. The model describes the foam bubble flow as well as the companion water flow  
392 during foam spurt. The maximum penetration depth by foam bubbles is predicted with a  
393 simplified micro stability model. The micro stability model results in good predictions when  
394 the foam bubbles are bigger than the pore throat in the sand while it underestimates the  
395 maximum penetration depth in other cases (for example, the tests with Sand 2 – FER 10 and  
396 Sand 3). Model results show good agreements when compared with experimental results during  
397 foam spurt.



398 The good agreement between the model and experimental results confirms the mechanisms  
399 during pure foam infiltration. Although further research is needed for sandy foam infiltration,  
400 which seems to be a more realistic situation in an EPB shield, pure foam infiltration is still a  
401 possible scenario during standstill as foam is injected through the cutter head to maintain the  
402 pressure from time to time. The injection will create a pure foam infiltration which is described  
403 in this paper.

404 Despite the fact that the model neglects the infiltration behavior after foam spurt, it provides a  
405 good way to quantify the flow behavior during foam infiltration. Further application of the  
406 model can be extended to the prediction of the pore pressures in front of the tunnel face, which  
407 could help better understand the pressure development and the foam-soil interaction in EPB  
408 shield tunneling. It is therefore recommended to engineers to incorporate the findings of this  
409 study in field use of foam for conditioning.

#### 410 **Acknowledgement**

411 The foaming agent used in this study was provided by the Institute for Tunnelling and  
412 Construction Management of Ruhr University Bochum. The help from Marius Schroeer and  
413 Sascha Freimann is highly appreciated. The soil water retention curve was measured at the  
414 Institute for Foundation Engineering, Soil and Rock Mechanics of Ruhr University Bochum,  
415 the assistance from Dr. Diethard Koenig is acknowledged. The first author would like to  
416 acknowledge the scholarship funded by China Scholarship Council and the CWO Mobility  
417 Fund of Ghent University.

#### 418 **Reference**

- 419 Anagnostou, G. & Kovári K., 1994. The face stability of slurry-shield-driven tunnels. TUST 9 (2): 165-174.  
420 Bezuijen, A., Schaminée, P.E.L. & Kleinjan, J.A., 1999. Additive testing for earth pressure balance shields.  
421 Proc. 12th Eur. Conf. on Soil Mech. and Geotech. Engrg., Amsterdam, Balkema, Rotterdam, pp. 1991–  
422 1996.  
423 Bezuijen, A., Pruiksmá, J.P. & Meerten, H.H., 2001. Pore pressures in front of tunnel, measurements,  
424 calculations and consequences for stability of tunnel face. Proc. Int. Symp. on Modern Tunneling  
425 Science and Techn. Kyoto.  
426 Bezuijen, A., 2002. The influence of soil permeability on the properties of a foam mixture in a TBM. 4th Int.  
427 Symp. on Geotechnical Aspects of Underground Construction in Soft Ground - IS Toulouse 2002.  
428 Bezuijen, A., 2012. Foam used during EPB tunnelling in saturated sand, parameters determining foam  
429 consumption. Proc. WTC, Bangkok.  
430 Bezuijen, A., Steeneken, S.P. et al., 2016. Monitoring and analysing pressures around a TBM. In: Proceeding

- 431 of the 13th International Conference Underground Construction, Prague.
- 432 Bezuijen, A. & Dias, T.G.S., 2017. EPB, chamber pressure dissipation during standstill. *EURO: TUN 2017*,  
433 Innsbruck University, Austria.
- 434 Broere, W. & van Tol A.F., 2000. Influence of Infiltration and Groundwater Flow on Tunnel Face Stability.  
435 Geotechnical Aspects of Underground Construction in Soft Ground, Tokyo, Japan, 2000, pp. 339-344.
- 436 Broere, W. & van Tol, A.F., 2001. Time-dependant infiltration and groundwater flow in a face stability  
437 analysis. In: Proceedings of International Symposium on Modern Tunnelling Science and Technology,  
438 Kyoto, pp. 629–634.
- 439 Budach, C., Thewes, M., 2015. Application ranges of EPB shields in coarse ground based on laboratory  
440 research. *Tunnelling and Underground Space Technology*, 50: 296-304.
- 441 Galli, M., 2016. Bochum Rheological characterisation of Earth-Pressure-Balance (EPB) support medium  
442 composed of non-cohesive soils and foam. Doctoral thesis. Ruhr-Universität, Germany.
- 443 Galli, M., Thewes, M., Freimann, S. & Schröer, M., 2021. Residual water content of excavated soil in EPB  
444 tunnelling. *TUST*, 114 (2021): 103991. doi.org/10.1016/j.tust.2021.103991
- 445 Lins, Y., 2009. Hydro-Mechanical Properties of Partially Saturated Sand. Doctoral thesis. Ruhr-Universität,  
446 Germany.
- 447 Maidl, U., 1995. Erweiterung der Einsatzbereiche der Erddruckschilde durch Bodenconditionierung mit  
448 Schaum. Doctoral Thesis, Ruhr-Universität Bochum, AG Lei-tungsbau und Leitungsinstandhaltung.
- 449 Merritt, A. S. & Mair, R. J. 2006. Mechanics of tunnelling machine screw conveyors: model tests.  
450 *Géotechnique* 56 (9): 605–615.
- 451 Min, F., Zhu, W., & Han, X., 2013. Filter cake formation for slurry shield tunneling in highly permeable sand.  
452 *TUST* 38(9): 423-430.
- 453 Peila, D., Oggeri, C. & Borio, L., 2009. Using the slump test to assess the behaviour of conditioned soil for  
454 EPB tunnelling. *Environmental & Engineering Geoscience XV* (3): 167–174.
- 455 Psomas, S., 2001. Properties of foam/sand mixtures for tunneling applications. Master's thesis, Oxford.
- 456 Quebaud, S., Sibai, M. & Henry, J.P., 1998. Use of chemical foam for improvements in drilling by earth  
457 pressure balance shields in granular soils. *TUST* 13 (2): 173-180.
- 458 Rossen. W.R. 1990. Theory of Mobilization Pressure Gradient of Flowing Foams in Porous Media. *Journal*  
459 *of Colloid and Interface Science* 136 (1): 1-16.
- 460 Rossen, W.R. & Gauglitz, P.A., 1990. Percolation Theory of Creation and Mobilization of Foams in Porous  
461 Media. *AIChE J.* 36: 1176–1188.
- 462 Talmon, A. M., Mastbergen, D. R., & Huisman, M., 2013. Invasion of pressurized clay suspensions into  
463 granular soil. *Journal of Porous Media* 16: 351-365.
- 464 Thewes, M., 2007. TBM Tunnelling Challenges—Redefining the State-of-the-Art, Keynote lecture at the  
465 2007 ITA World Tunnel Congress, Prague, *Tunel*. pp. 13–21.
- 466 Thewes, M., Budach, C. & Bezuijen, A., 2012. Foam conditioning in EPB tunneling. *Geotechnical Aspects*  
467 *of Underground Construction in Soft Ground*. London.
- 468 USACE (2000) Engineering Manual 1110-2-1913, Design and Construction of Levees.  
469 [https://www.publications.usace.army.mil/Portals/76/Publications/EngineerManuals/EM\\_1110-2-](https://www.publications.usace.army.mil/Portals/76/Publications/EngineerManuals/EM_1110-2-1913.pdf)  
470 [1913.pdf](https://www.publications.usace.army.mil/Portals/76/Publications/EngineerManuals/EM_1110-2-1913.pdf).
- 471 van Genuchten, M., 1980. A closed-form equation for predicting the hydraulic conductivity of unsaturated  
472 soils. *Soil Sci. Soc. Am. J.* 44: 892–898.
- 473 Vargaftik, N.M., Volkov, B.N. & Voljak, L.D., 1983. International Tables of the Surface Tension of Water.  
474 *Journal of Physical and Chemical Reference Data*. 12 (3): 817-820.

- 475 Wang, S., Hu, Q., Wang, H., Thewes, M., Ge, L., Yang, J. & Liu, P., 2021. Permeability Characteristics of  
476 Poorly Graded Sand Conditioned with Foam in Different Conditioning States. *Journal of Testing and*  
477 *Evaluation*. Published ahead of print, September 01, 2021. <https://doi.org/10.1520/JTE20190539>
- 478 Wei, Y., Yang, Y. & Qiu, T., 2019. Effects of Soil Conditioning on Tool Wear for Earth Pressure Balance  
479 Shield Tunneling in Sandy Gravel Based on Laboratory Test. *Journal of Testing and Evaluation* 49.  
480 <https://doi.org/10.1520/JTE20180851>.
- 481 Wu, Y., Nazem, A., Meng, F. & Mooney, M.A., 2020. Experimental study on the stability of foam-conditioned  
482 sand under pressure in the EPBM chamber. *TUST* 106 (2020): 103590.
- 483 Xu, T., 2018. Penetration and excess pore water pressures in front of a TBM, Experiments, Mechanisms and  
484 Computational models. Doctoral thesis. Ghent University, Belgium.
- 485 Xu, T. & Bezuijen, A., 2019. Bentonite slurry infiltration into sand: filter cake formation under various  
486 conditions. *Géotechnique*. <https://doi.org/10.1680/jgeot.18.P.094>
- 487 Xu, T., Bezuijen, A. & Thewes, M., 2021. Pressure infiltration characteristics of foam for EPB shield  
488 tunnelling in saturated sand – part 1: ‘clean’ foam. *Géotechnique*. <https://doi.org/10.1680/jgeot.19.P.187>
- 489 Xu, T., Bezuijen, A. & Thewes, M., 2021. Pressure-infiltration characteristics of foam for EPB shield  
490 tunnelling in saturated sand – part 2: soil–foam mixture. *Géotechnique*.  
491 <https://doi.org/10.1680/jgeot.19.P.188>
- 492 Yin, T., Zhang, Z., Huang, X., Shire, T. & Hanley, K.J., 2021. On the morphology and pressure-filtration  
493 characteristics of filter cake formation: Insight from coupled CFD–DEM simulations. *TUST* 111 (2021):  
494 103856. <https://doi.org/10.1016/j.tust.2021.103856>
- 495 Yu, H., Mooney, M. & Bezuijen, A., 2020. A simplified excavation chamber pressure model for EPBM  
496 tunneling. *Tunnelling and Underground Space Technology* 103 (2020): 103457.  
497 <https://doi.org/10.1016/j.tust.2020.103457>
- 498 Zheng, D., Bezuijen, A. & Thewes, M., 2021. An experimental study on foam infiltration into saturated sand  
499 and its consequence for EPB shield tunneling. *TUST*, 111 (2021): 103878.  
500 <https://doi.org/10.1016/j.tust.2021.103878>
- 501 Zizka, Z., Schoesser, B., Thewes, M., Schanz, T., 2018. Slurry shield tunneling: New methodology for  
502 prediction of increased pore pressures resulting from slurry infiltration at the tunnel face under cyclic  
503 excavation processes. *International Journal of Civil Engineering* 7: 113–130.

504 **Tables**

505

Table 1. Values of  $D_{15}/d_{85}$  found in different combinations

	Sand 1	Sand 2	Sand 3
Foam 1	0.29	0.28	0.44
Foam 2	0.26	0.28	0.56
Foam 3	0.21	0.22	0.44

506

Table 2. Parameters found in experiments and model simulation

	FER	$d_{10,s}$ (m)	$d_{10,p}$ (m)	$d_{10,f}$ (m)	Volume factor ( $\alpha$ )	$L$ (m)		
						Micro stability model	Model fit	Experiment*
Sand 1	10	$1.2 \times 10^{-4}$	$3.5 \times 10^{-5}$	$1.8 \times 10^{-4}$	4	0.016	0.017	$0 < L < 0.02$
	15			$2.4 \times 10^{-4}$	8	0.033	0.036	$0.04 < L < 0.06$
	20			$2.6 \times 10^{-4}$	11	0.045	0.045	$0.04 < L < 0.06$
Sand 2	10	$1.3 \times 10^{-4}$	$5.5 \times 10^{-5}$	$1.8 \times 10^{-4}$	3	0.026	0.06	$0.04 < L < 0.06$
	15			$2.4 \times 10^{-4}$	7	0.061	0.061	$0.04 < L < 0.06$
	20			$2.6 \times 10^{-4}$	8	0.069	0.069	$0.06 < L < 0.1$
Sand 3	10	$2.6 \times 10^{-4}$	$7.0 \times 10^{-5}$	$1.8 \times 10^{-4}$	1	0.017	0.09	$0.06 < L < 0.1$
	15			$2.4 \times 10^{-4}$	1	0.016	0.09	$0.06 < L < 0.1$
	20			$2.6 \times 10^{-4}$	1	0.023	0.09	$0.06 < L < 0.1$

507 \* Experimental data from Zheng et al. (2021)

508

Table 3. Measured capillary rise and calculated surface tension

	$h$ (m)	$\gamma$ (N/m)
water	0.1982	0.0728
Foaming liquid	0.0583	0.0214

509

Table 4. Permeabilities found in experiments and model simulation

	$k_s$ ( $\times 10^{-4}$ m/s)		FER 10 ( $\times 10^{-8}$ m/s)			FER 15 ( $\times 10^{-8}$ m/s)			FER 20 ( $\times 10^{-8}$ m/s)		
	$k_s$ (Fig. 8)	$k_s^*$	$k_{fw}$ (Fig.8)	$k_{fw}$	$k_f$	$k_{fw}$ (Fig.8)	$k_{fw}$	$k_f$	$k_{fw}$ (Fig.8)	$k_{fw}$	$k_f$
Sand 1	1.0	1.5	2	2.6	8	2	1.8	4	2	1.5	3
Sand 2	4.0	5.0	2	4.3	42	3	4	30	2	2.3	25
Sand 3	6.0	6.0	15	5.8	80	10	3.6	45	5	2.3	40

510 \* Experimental data from Zheng et al. (2021)

511 **Figure Captions List**

- 512 Figure 1. Sketch of mechanisms during foam infiltration into saturated sand.
- 513 Figure 2. Definition sketch of penetration zone and piezometric head during foam penetration  
514 (set-up described in Zheng et al., 2021).
- 515 Figure 3. Experimental results of the water discharge against square root of time (data from  
516 Zheng et al., 2021)
- 517 Figure 4. Grain / Bubble size distribution curves of the sands / foams (data from Zheng et al.,  
518 2021).
- 519 Figure 5. Principle of the bubble deformation during foam penetration of a pore throat.
- 520 Figure 6. Possible conditions when foam bubbles penetrate through pore throats.
- 521 Figure 7. Measured water retention curves with van Genuchten model fit (a) and the distribution  
522 curve of the throat diameter for each sand (b).
- 523 Figure 8. Calculated permeabilities and specific discharge as a function of infiltration depth  
524 (data from Zheng et al., 2021).
- 525 Figure 9. Penetration depth against time from experiments compared to model results.
- 526 Figure 10. Permeabilities used in model simulation.





















































